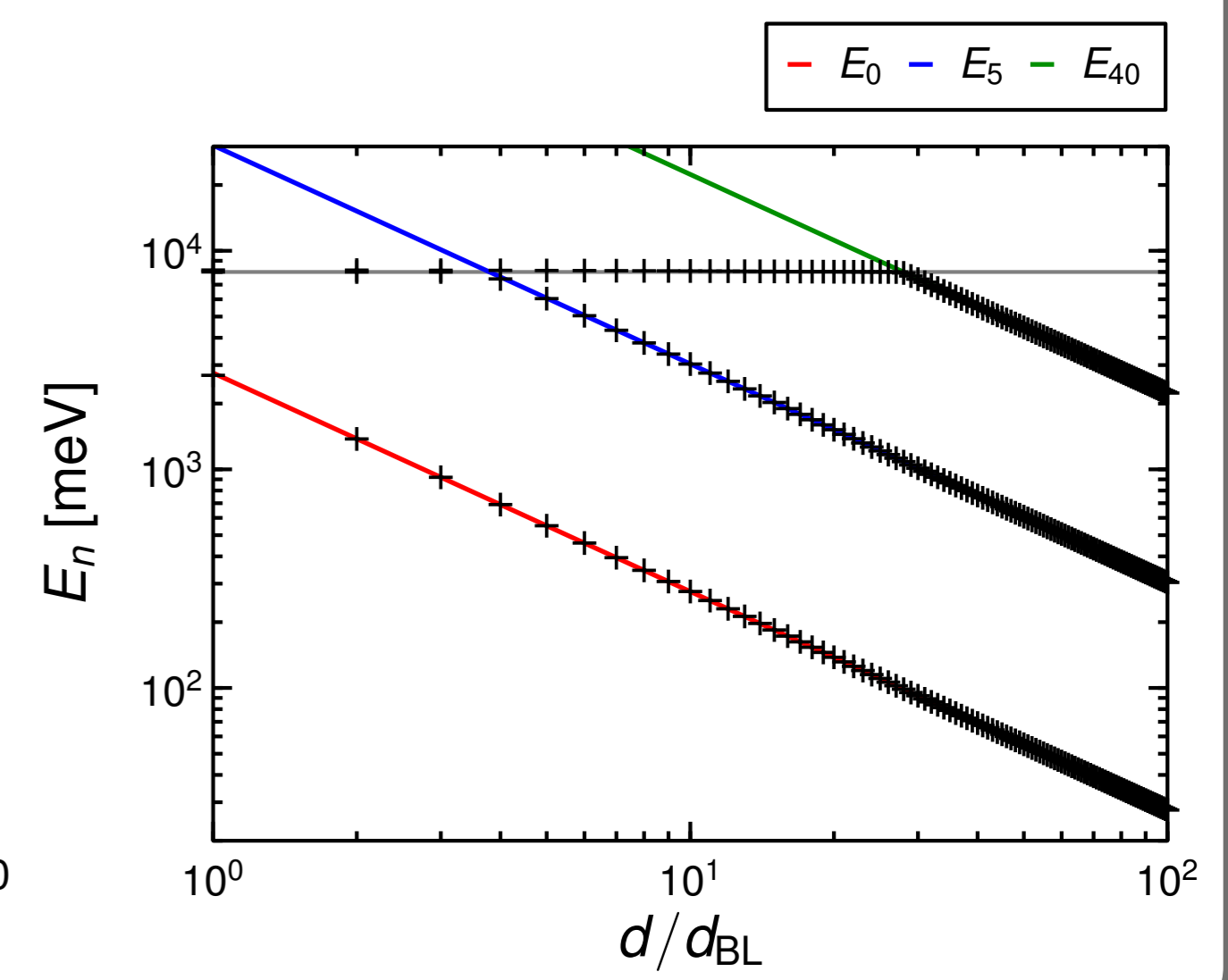
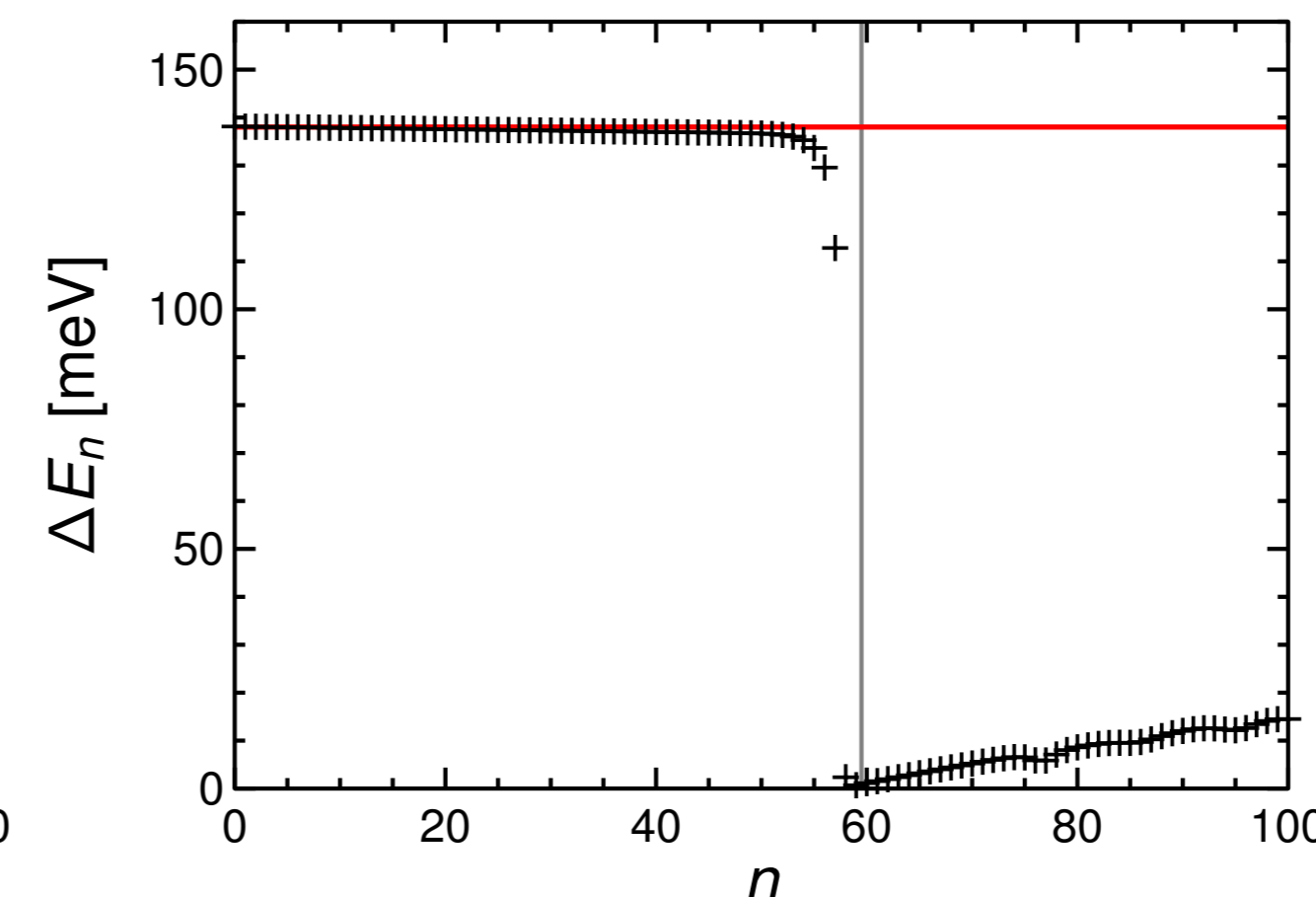
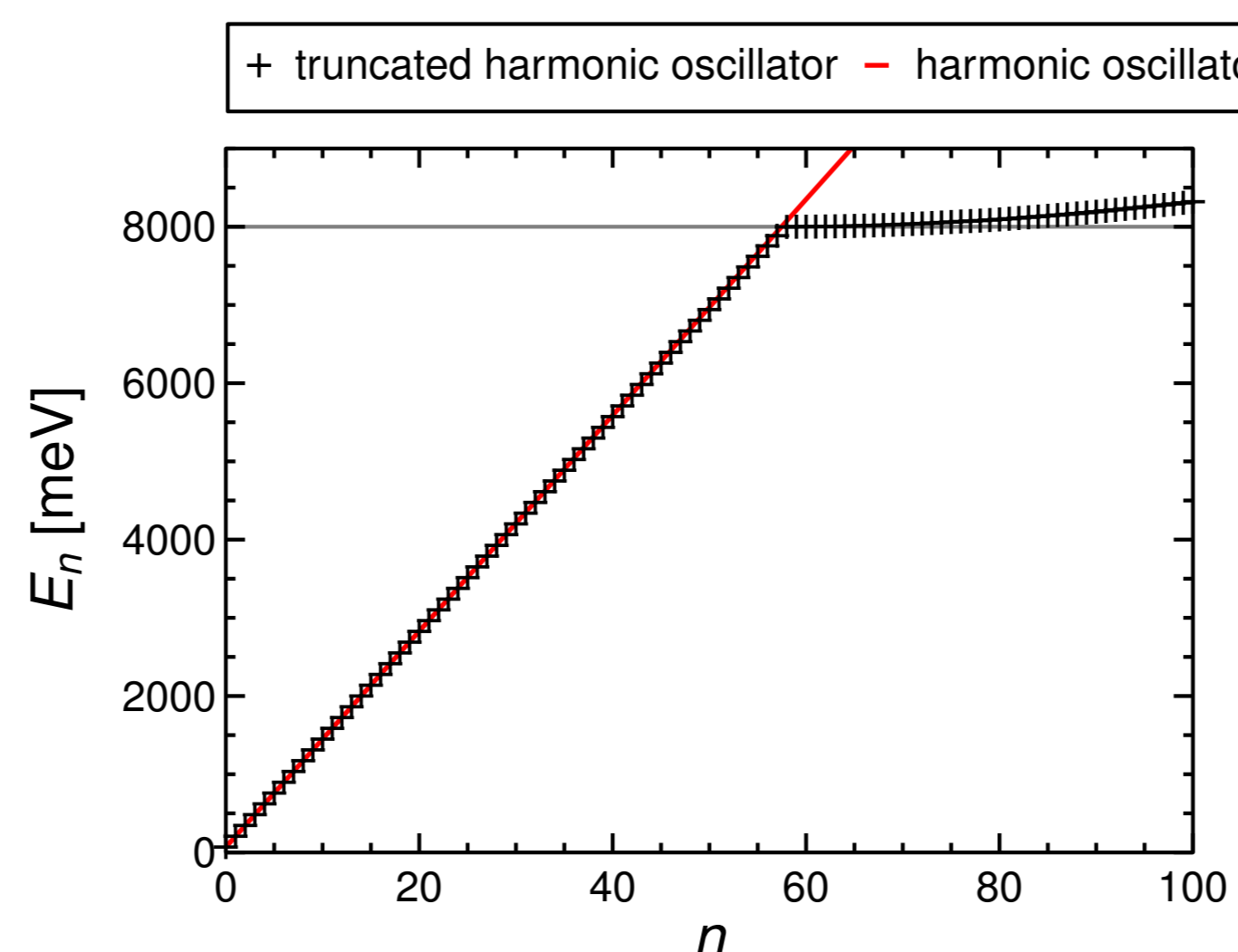
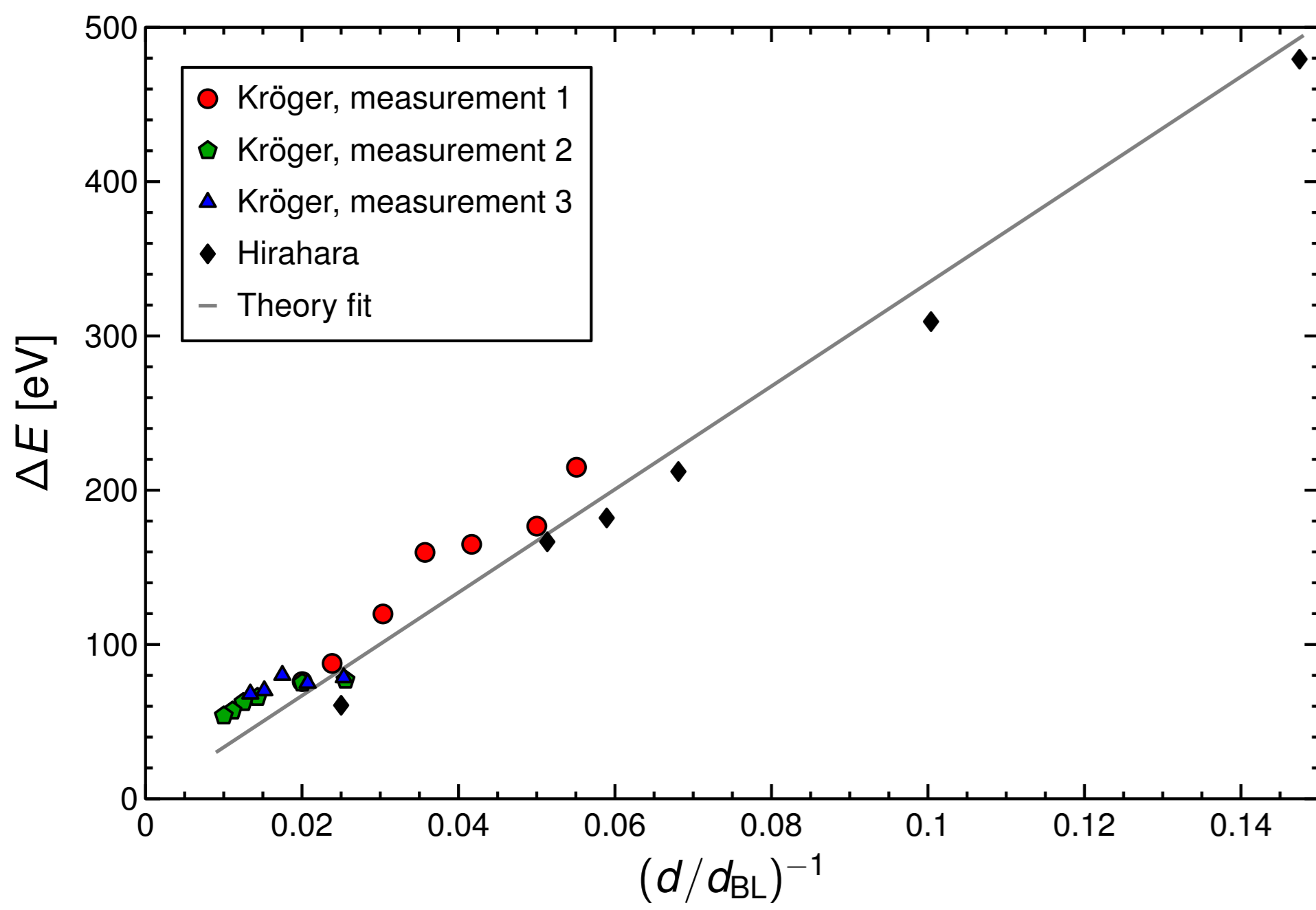




## Experimental results and harmonic oscillator model

- [1] T. Hirahara et al.: *Quantum well states in ultrathin Bi films: Angle-resolved photoemission spectroscopy and first-principles calculations study*, Physical Review B 75 (2007), 035422
- [2] P. Kröger, C. Tegenkamp et al.: *Controlling conductivity by quantum well states in ultrathin Bi(111) films*, Physical Review B 97 (2018), 045403



- Quantum well states in Bi films: bandgap  $E_g$  measured,  $1/d$ -dependence found ( $d$  ... film thickness)
- Typical models: Square well potential:  $\Delta E \sim 1/d^2$ , Quadratic potential:  $\Delta E \sim 1/d$

- Simple model: truncated harmonic oscillator  $V(x) = E_g \begin{cases} 1 & |x| \geq d/2 \\ (2x/d)^2 & |x| \leq d/2 \end{cases}$   
 $\Delta E = \frac{2\hbar}{d} \sqrt{\frac{2E_g}{m^*}}$
- Example:  $E_g/m^* = 8 \text{ eV}/m_e$ ,  $d = 16 \text{ nm}$  (fig. below)

## Approaches for calculating potentials with equidistant spectra

### Shift operator approach

- [3] Dubov et al.: *Equidistant spectra of anharmonic oscillators*, Journal of Experimental and Theoretical Physics 75 (1992), p. 446–451
- [4] S. Dubov et al.: *Equidistant spectra of anharmonic oscillators*, Chaos 4 (1994), p. 47–53

### Factorization method

- [5] B. Mielnik: *Factorization method and new potentials with the oscillator spectrum*, Journal of Mathematical Physics 25 (1984), p. 3387–3389
- [6] D. J. Fernández C et al.: *A simple generation of exactly solvable anharmonic oscillators*, Physics Letters A 244 (1998), p. 309–316

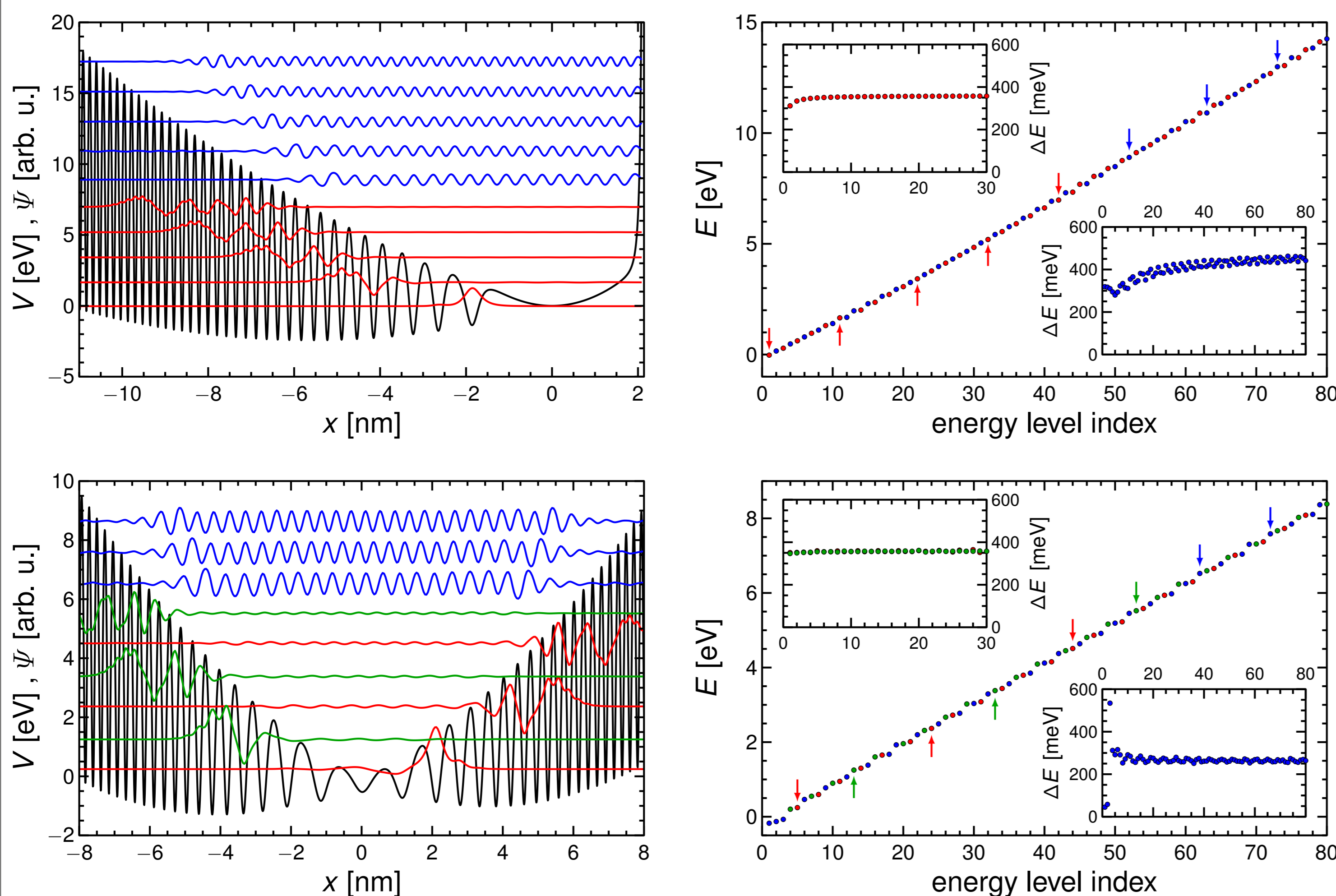
## Conclusion

- Energy levels of thin films cannot provide all information about the underlying confinement potentials
- Many different oscillating and diverging potential shapes with equidistant energy levels are possible
- Different spectra are constructable with factorization method choosing appropriate  $\epsilon_{k_i}$
- Different potential shapes for given spectrum are adjustable by choosing appropriate  $A_{k_i}$

## Shift operator approach [3,4]

- Solving  $[\mathcal{H}, \mathcal{L}] = \mathcal{L}$ ,  $\mathcal{L}$  ... shift operator (generalized creation operator)
- Expansion  $\mathcal{L} = \sum_{k=0}^K \alpha_k(\xi) (i\mathcal{P})^k$ ,  $\mathcal{P}$  ... momentum operator
- Solving  $k+1$  differential equations for unknowns  $\alpha_k(\xi)$  and unknown potential  $U(\xi)$
- 1st order:  $\mathcal{L} = \alpha_0(\xi) + \alpha_1(\xi) i\mathcal{P} \rightarrow$  creation operator  $\rightarrow$  harmonic potential  $U(\xi) = \frac{1}{2}\xi^2$
- 2nd order: isotonic potential  $U(\xi) = \frac{1}{8}\xi^2 + \frac{A}{\xi^2}$
- 3rd order: Solving  $\frac{3}{2}[W(\xi)^2]'' - \frac{1}{4}W''''(\xi) - \xi^2 W''(\xi) + 3\xi W'(\xi) = 0$  with  $U(\xi) = \frac{1}{2}\xi^2 + W(\xi)$
- Darboux transformation yields some solutions:  
 $U_m(\xi) = -\frac{1}{2}\xi^2 + \left(\frac{P'_m(\xi)}{P_m(\xi)} + \xi\right)^2$ ,  $P_{2m+\delta} = \sum_{k=0}^m \frac{4^k}{(m-k)!(2k+\delta)!} \xi^{2k+\delta}$ ,  $\delta \in \{0, 1\}$

### Exemplary numeric solutions:



## Factorization method [5,6]

- Operator  $a^\dagger = \frac{1}{\sqrt{2}} \left[ -\frac{d}{d\xi} + \alpha(\xi) \right]$  transforms  $\mathcal{H}_1$  into  $\mathcal{H}_2$ :  $\mathcal{H}_2 a^\dagger = a^\dagger \mathcal{H}_1$  with  $\mathcal{H}_i = -\frac{1}{2} \frac{d^2}{d\xi^2} + V_i(\xi)$
- $V_2(\xi) = V_1(\xi) - \alpha'(\xi) + \alpha(\xi)^2 = 2[V_1(\xi) - \epsilon]$
- $V_2(\xi)$  has same spectrum as  $V_1(\xi)$ , except an additional energy level at  $\epsilon$
- Analytically solvable for  $V_1(\xi) = \frac{1}{2}\xi^2$  for  $\epsilon = -\left(k + \frac{1}{2}\right) \rightarrow \alpha_k(\xi), V_k(\xi)$   
 $\alpha_k(\xi) = \xi + \frac{d}{d\xi} \log \left\{ [1 + A_k \text{erf}(\xi)] h_k(\xi) + A_k \frac{2}{\sqrt{\pi}} e^{-\xi^2} p_k(\xi) \right\}$  with  $h_k(\xi) = (-i)^k H_k(\xi)$ ,  
 $p_{2k+\delta}(\xi) = \sum_{m=0}^k \frac{2^m (2k-m-\delta)!}{(2k-2m-\delta)!} h_{2k-2m-\delta}(\xi)$ ,  $\delta \in \{0, 1\}$ ,  $A_k$  ... constant
- Multiple transformations with different  $\epsilon_{k_i}$ : additional energy levels at every  $\epsilon_{k_i}$   
 $V_{k_1 \dots k_n}(\xi) = V_{k_1 \dots k_{n-1}}(\xi) - \alpha'_{k_1 \dots k_n}(\xi)$ ,  $\alpha_{k_1 \dots k_n}(\xi) = -\alpha_{k_1 \dots k_{n-1}}(\xi) - 2 \frac{k_n - k_{n-1}}{\alpha_{k_1 \dots k_{n-1}}(\xi) - \alpha_{k_1 \dots k_{n-2} k_n}(\xi)}$

### Exemplary solutions:

