

# Complex dynamics made simple: colloidal dynamics

Paolo Sibani

Chemnitz, June 2010

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- Age: time elapsed from the initial quench
- Time homogeneous dynamics: e.g. diffusion, no changes of mesoscopic rates with time.
- Time inhomogeneous dynamics: e.g. aging: rate of significant events goes down with age.



# Metastability and coarse-graining in energy landscapes

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# Metastability and coarse-graining in energy landscapes

- Many basins of attraction. E.g. energy minima separated by energy barriers.
- On long enough time scales neglect the details and look at the dynamics at the level of attractor changes.
- The resulting dynamics may be simple: e.g. diffusion in sinusoidal potential
- Or complex: e.g. spin glass dynamics, where averages have power-law or logarithmic dependences.

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- Hard sphere colloidal systems-only short range interactions
- Entropy driven dynamics.
- Clusters of highly correlated groups of particles (spatial coarse graining)
- Dynamics described in terms of the probability per unit of time that a cluster survives the random forces from the rest of the system.



# What does 'trivial' or 'simple' mean?

- Trivial dynamics  $\stackrel{\text{def}}{=}$  time homogeneous (but perhaps not stationary)







# The simplest stochastic process

- Add random and independent 'position' increments

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- Use linearity of the variance
- $MSD \propto t$  : diffusion

## Record dynamics

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- record sized fluctuations
- in the stationary white noise impinging on the system (e.g. thermal noise)
- drive the aging dynamics in marginally stable (glassy) systems
- no. of records is a **Poisson process in logarithmic time**
- the dynamics appears simple in logarithmic time

# The rest of this talk

- Brief overview of results for *thermally activated* dynamics

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- Experimental<sup>1</sup> colloidal data re-analyzed.

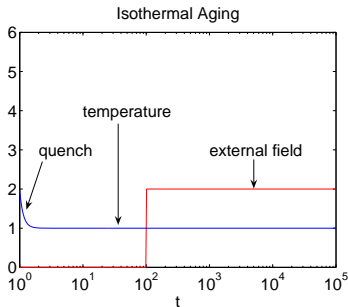
## The rest of this talk

- Brief overview of results for *thermally activated* dynamics
- Experimental<sup>1</sup> colloidal data re-analyzed.
- A cluster model for colloidal behavior numerically analyzed.





# Aging set-up



The age  $t$  starts at the 'initial quench', and a field is (possibly) switched on at  $t = t_w = 100$ . In a colloid, the system is centrifuged to a high density (quench). No fields are applied.

# Edwards Anderson spin glass

The Edwards Anderson spin-glass is an Ising spin model where interacting spins  $\sigma_i = \pm 1$  are placed on a lattice. The Hamiltonian is

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- For nearest neighbors  $J_{ij}$  are Gaussian standard random variables, independent for  $i < j$
- The thermal equilibrium properties of the model are *complex*
- The time evolution, as given by a MC algorithm (Metropolis acceptance rule or equivalent) is *complex* and very similar to experimental data

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- $\mathcal{H} = - \sum_{Plaq.} \sigma_i \sigma_j \sigma_k \sigma_l$
- The thermal equilibrium properties of the model are *trivial*
- The time evolution is nevertheless complex, featuring metastability and aging.



The Restricted Occupancy Model is a lattice model describing vortex creep in type II superconductors in terms of the number  $n_i =$  of vortices on site  $i$ .

- the energy includes repulsive interactions between vortex lines on neighbor sites

see L. P. Oliveira et al. *Phys. Rev. B* 104526 (2005) and references therein.

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- pinning to random sites

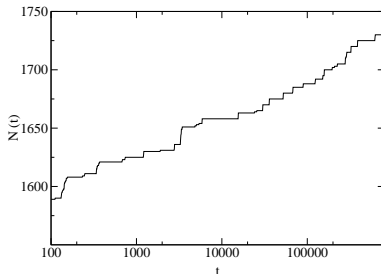
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- the energy includes repulsive interactions between vortex lines on neighbor sites
- pinning to random sites
- the configuration is updated with Metropolis dynamics

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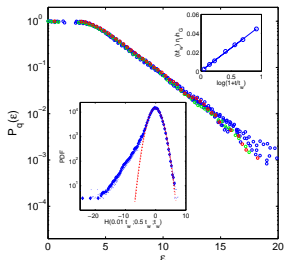
# ROM model intermittency



The time variation of the total number of vortices  $N(t)$  on the system for a single realization of the pinning potential and the thermal noise in a  $8 \times 8 \times 8$  lattice for  $T = 0.1$ .

L. P. Oliveira et al. Phys. Rev. B 104526 (2005)

# Spin-glass, heat flow intermittency

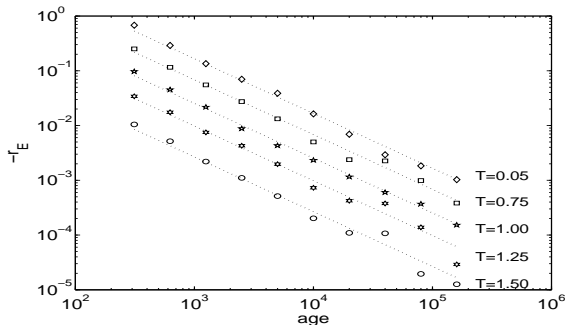


Heat transfer PDF for a spin glass model (PS & H J Jensen, Europhysics Lett. 2005)

Heat

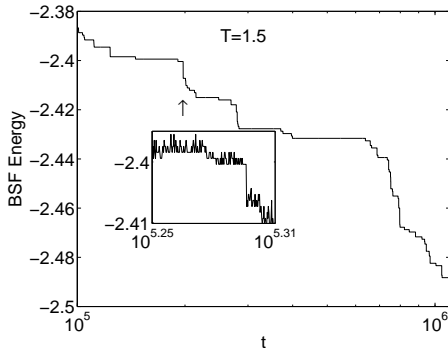
transfer  $H$  over small time  $\delta t$  in the E-A spin glass model has a Gaussian part and an intermittent tail. Six different ages are considered with  $\delta t/t_w = .01$ .

# Heat flow rate, p-spin model



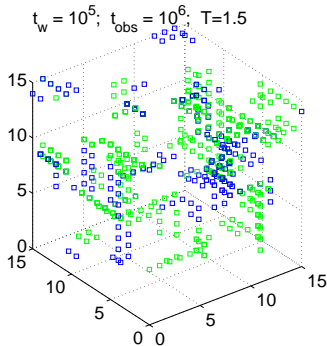
The average rate of energy flow is plotted versus the age for the temperatures shown. The full line has the form  $y = C(T)t_w^{-1}$ .

# p-spin model, intermittent energy decay



S. Christiansen  
& PS, *New J. of  
Physics*, to appear

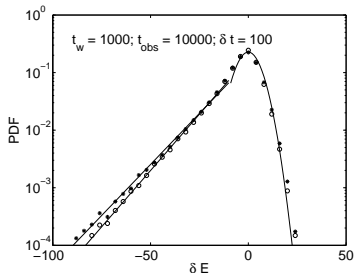
# p-spin model, quakes in real space



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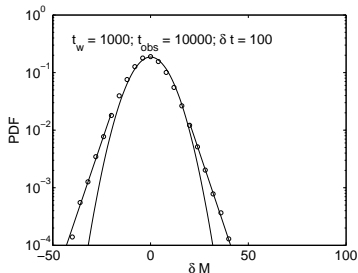


# p-spin model, heat flow



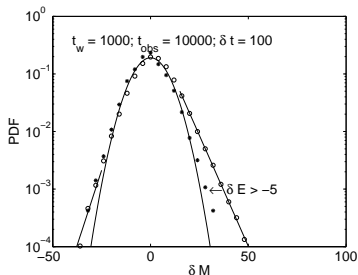
The PDF of the heat exchanged between system and thermal bath over a time  $\delta t = 100$ .  $T = 1.5$ .

# p-spin model, magnetic fluctuations $H = 0$



The PDF of the spontaneous magnetic fluctuations over a time  $\delta t = 100$  and  $T = 1.5$ .

# p-spin model, magnetic fluctuations $H = 0.3$



PDF of the spontaneous magnetic fluctuations over a time  $\delta t = 100$ .  
 $T = 1.5$ . The

magnetic response is SUBORDINATED to the quakes

## Take-home message from thermally activated aging

- The evolution of aggregate variables (i.e. the energy and magnetization) is controlled by extremely rare and irreversible events **quakes**

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- The evolution of aggregate variables (i.e. the energy and magnetization) is controlled by extremely rare and irreversible events **quakes**
- Reversible fluctuation of zero average with Gaussian PDF's describe the dynamics in quasi-equilibrium.

- Tracking data obtained by confocal microscopy

Courtland and Weeks, J. Phys.: Condens. Matter 15, S359, (2003)

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- dense colloids  $\rho > 0.62$  sub-diffusive behavior, aging
- 'glass formers'  $\rho < 0.62$  diffusive behavior, time homogeneous

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## Experimental procedure

- Centrifuge the sample to desired density

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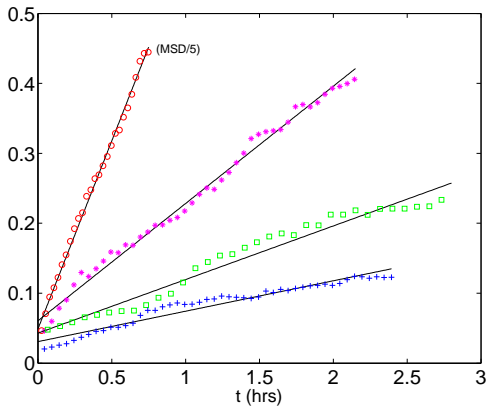
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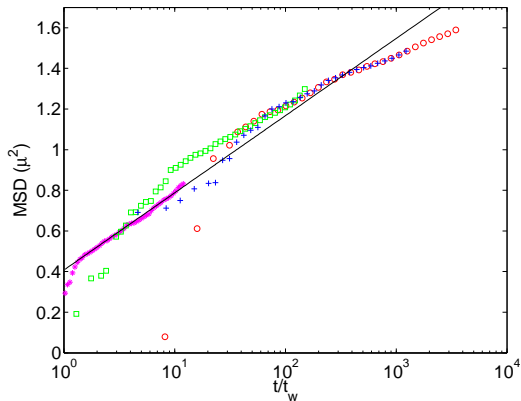
- Centrifuge the sample to desired density
- Stir the sample
- Wait (age  $t_w$  calculated from the end of the stirring phase)



# MSD vs time. Glass-former



# MSD vs time. Dense





- Partition the system volume in a set of subvolumes

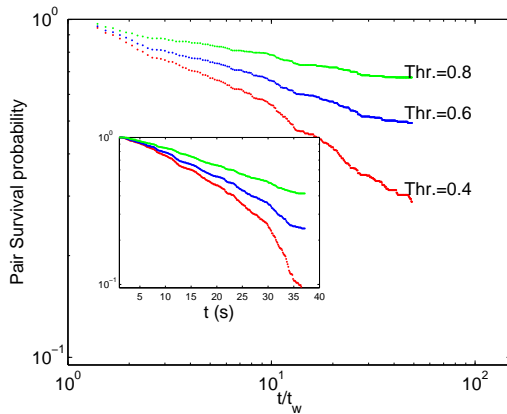
# Persistence analysis

- Partition the system volume in a set of subvolumes
- Pick for each sub-volume a pair of colloidal particle which are initially 'in touch'

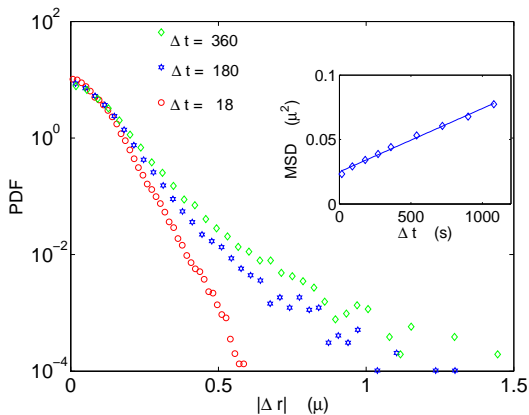
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- Calculate the fraction of pairs which survive at time  $t$ .

# Persistence curves

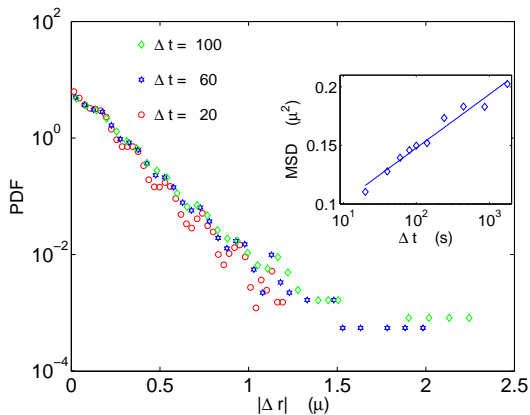


# Intermittency curves



'Glass former'

# Intermittency curves



# Conclusion from colloid data analysis

- the MSD is linear in  $t$  (glass formers) and in  $\ln(t) - \ln(t_w)$  (dense)



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- Transformation  $t \rightarrow \ln(t/t_w)$  trivializes the dynamics of dense colloids
- Highly intermittent and correlated motion

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- Particle density is constant

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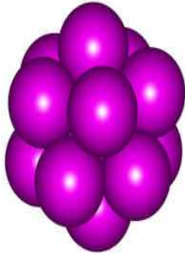
- Particle density is constant
- Energy density is constant



## Not much changes in a colloid

- Particle density is constant
- Energy density is constant
- Particle displacement over observation time of the order of particle radius
- Particle motion is jerky and spatially correlated

# A cluster model



particles in a cluster moves  
in fully correlated fashion  
on average zero CM displacement  
as long as the cluster persists



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- for all  $k$ :  $P_k(h \rightarrow 0)$  for  $h \rightarrow \infty$

# The algorithm

- Markov chain,  $L$  clusters

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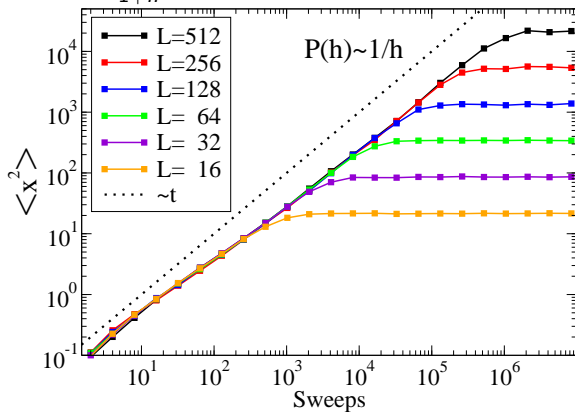
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# The algorithm

- Markov chain,  $L$  clusters
- pick a cluster at random
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- iff cluster is destroyed:
  - partition its particles randomly in two groups
  - each sub-group joins neighboring cluster
  - particles move one random unit step on the lattice

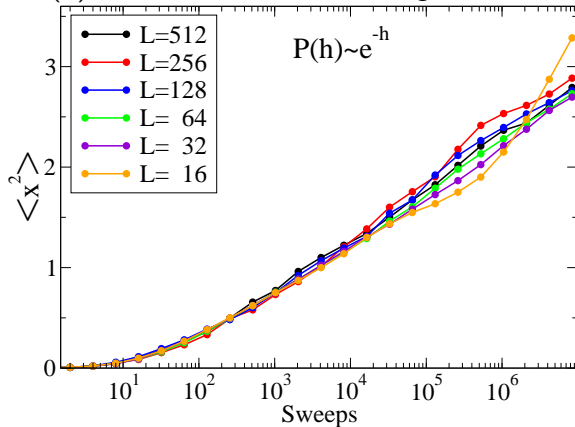
# Diffusive behavior I

$$P_1(h) = \frac{1}{1+h} \quad \text{'glass former', diffusive}$$



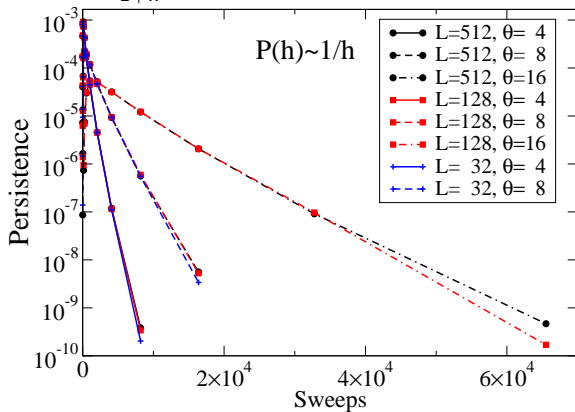
# Diffusive behavior II

$P_{\infty}(h) = e^{-h}$  dense colloid, log diffusive



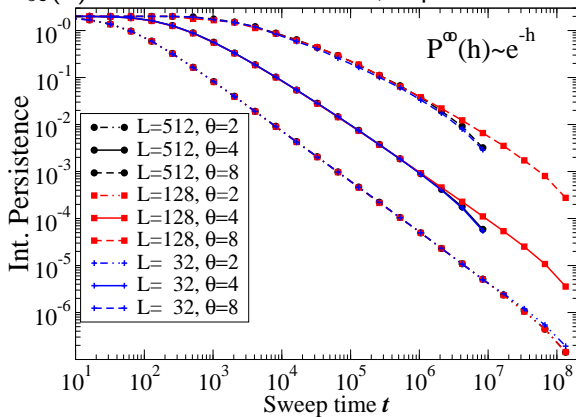
# Persistence I

$$P_1(h) = \frac{1}{1+h} \quad \text{'glass former', exponential decay}$$



# Persistence II

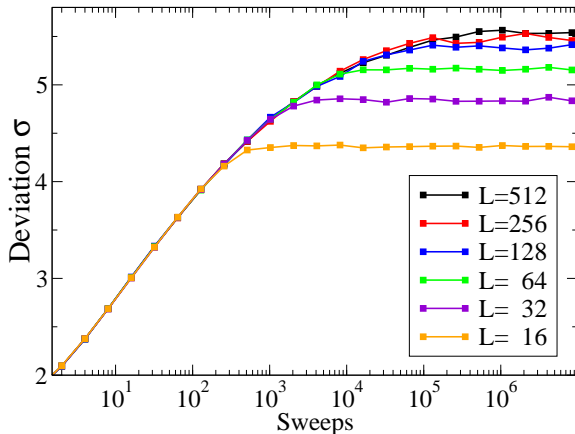
$P_{\infty}(h) = e^{-h}$  dense colloid, exponential decay in  $\log t$



# Spatial heterogeneity I

$$P_1(h) = \frac{1}{1+h} \quad \text{'glass former', approaches stationary state}$$

$P(h) \sim 1/h$

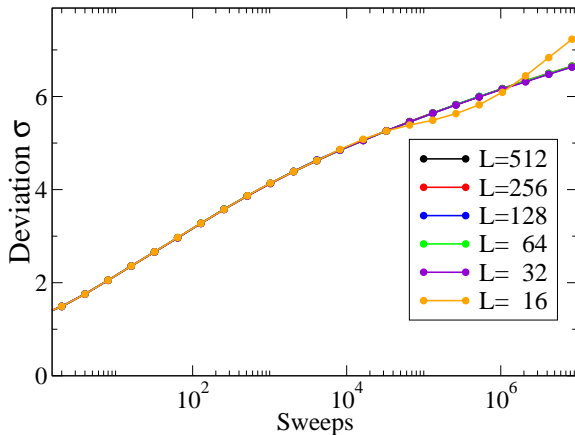




# Spatial heterogeneity II

$P_{\infty}(h) = e^{-h}$  dense colloid, equilibrium is out of reach

$P(h) \sim e^{-h}$



# Quakes in colloidal model

- Collapses of clusters of different sizes are widely separated in time
- All clusters of size  $< \sigma(t)$  only provide background of mobile particles
- A quake on a time scale  $t$  corresponds to the collapse of the smallest cluster  $> \sigma$

# Quakes in thermally activated aging

- Energy fluctuations which are negative

# Quakes in thermally activated aging

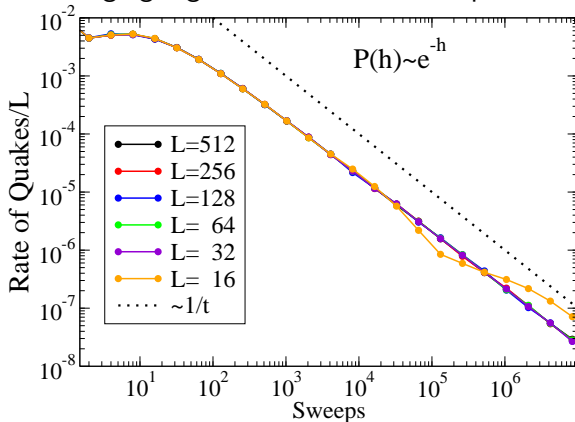
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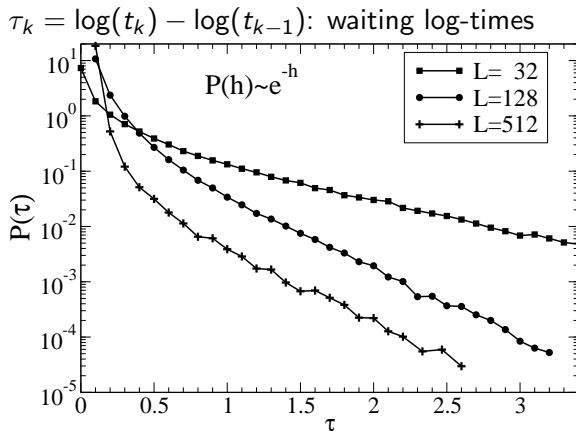
- Energy fluctuations which are negative
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- Size  $-\delta e/T \gg \log t$

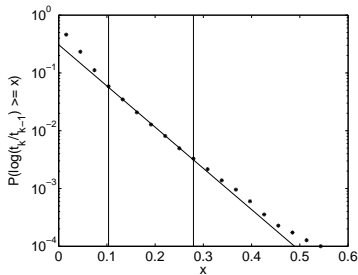
# Temporal statistics of quakes

In the aging regime, each cluster collapse is called a 'quake'



# Temporal statistics of 'quakes' II





The distribution of the logarithmic differences  $\log(t_k) - \log(t_{k-1})$  is approximately exponential.



# The 'log' Poisson distribution

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$$P(n, t_1, t_2) = \frac{\mu^n}{n!} \exp(-\mu) \quad \mu(t_1, t_2) = \alpha \log(t_2/t_1) \quad (1)$$

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- The statistics applies to e.g. spin-glasses, colloids, evolutionary dynamics etc.
- Where can it possibly come from?



# Relation to hierarchies and trees?

- Two different ways to describe the same generic situation

( Fischer Hoffmann & Sibani, PRE 2008 )

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- Record dynamics needs an underlining hierarchical structure

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# Relation to hierarchies and trees?

- Two different ways to describe the same generic situation
- Record dynamics needs an underlining hierarchical structure
- work with Andreas Fischer and KH elaborates on this point

( Fischer Hoffmann & Sibani, PRE 2008 )







# Acknowledgments

- Colloid analysis in collaboration with Stefan Boettcher, Emory University, GA, USA. Thanks to Eric Weeks Emory University, for kindly providing the colloidal data.
- The p-model data shown are part of a (former) student, Simon Christiansen's master thesis
- Thanks to Jesper Dall, Christian Schön Karl Heinz Hoffmann & Henrik J. Jensen for inputs and discussions over the years.