

DELONE MEASURES OF FINITE LOCAL COMPLEXITY AND
APPLICATIONS TO SPECTRAL THEORY OF
ONE-DIMENSIONAL CONTINUUM MODELS OF
QUASICRYSTALS

STEFFEN KLASSERT

Fakultät für Mathematik
TU Chemnitz
D - 09107 Chemnitz, Germany

DANIEL LENZ

Mathematisches Institut, Friedrich-Schiller Universität
Ernst-Abbe-Platz 2
D - 07743 Jena, Germany

PETER STOLLMANN

Fakultät für Mathematik
TU Chemnitz
D - 09107 Chemnitz, Germany

ABSTRACT. We study measures on the real line and present various versions of what it means for such a measure to take only finitely many values. We then study perturbations of the Laplacian by such measures. Using Kotani-Remling theory, we show that the resulting operators have empty absolutely continuous spectrum if the measures are not periodic. When combined with Gordon type arguments this allows us to prove purely singular continuous spectrum for some continuum models of quasicrystals.

Dedicated to Hajo Leschke on the occasion of his 65th birthday

Introduction. Quasicrystals are interesting objects, from a physical and a mathematical point of view. For the physics we refer to the pioneering work of Shechtman, Blech, Gratias and Cahn [34], published some 25 years ago, as well as to [4, 13, 33] and the references therein. The mathematical interest has a number of sources. Here, we are mainly concerned with the fact that quantum mechanical models of quasicrystal exhibit a number of “strange phenomena”: there is a tendency that the respective Hamiltonians have Cantor sets as spectra, with absence of absolutely continuous spectrum as well as absence of point spectrum. Moreover, anomalous transport is typical for these operators. Very roughly speaking, this can be motivated by the fact that quasicrystals are aperiodically ordered. They are, in a sense, close to periodic structures, which explains the absence of point spectrum. On the other hand, they can also be thought of as random operators, which explains the absence of absolutely continuous spectrum.

2000 *Mathematics Subject Classification.* Primary:81Q10, 35J10; Secondary: 52C32, 82D30.
Key words and phrases. Quasicrystals, Delone sets, absolutely continuous spectrum.

Starting with [2, 38, 39], the general paradigm of singular continuous Cantor spectrum has been verified in the discrete one-dimensional setting for various examples, see e.g. the surveys [7, 22] and references therein. The only results for multidimensional models we are aware of are from [25] and the only results in the one-dimensional continuum setting are from [15]. Here, we present a considerable extension of the latter results. In the discrete case, absence of absolutely continuous is a consequence of Kotani theory, [16, 17]. Absence of eigenvalues can be shown with Gordon-type arguments, see [7] for a survey and a list of relevant references.

It is quite reasonable to try similar strategies for continuum models; one part of that program, the respective Gordon-type arguments have already been established in [9] (see [32] for recent extensions as well). So the main point is to prove absence of absolutely continuous spectrum by an appropriate version of Kotani's theory. This had been done in the thesis of one of us, [15] within the framework of Kotani theory for random operators. In the meantime a deterministic version of Kotani theory has been developed by Remling [31]. In the present work we rely upon Remling's oracle theorem, that allows a purely deterministic approach. The outcome is Theorem 4.1 below that can also be applied in the random setting of measure dynamical systems.

We start with an account on Schrödinger operators with measure potentials since we adopt the generality of [31] in the following section. There, we will also state the above mentioned oracle theorem, our main tool to prove absence of absolutely continuous spectrum. We give a little discussion of finite local complexity for point sets in Section 2. Then, we present several different versions of a concept of finite local complexity in the context of potentials in Section 3 and study their relationship. This study may be of independent interest. The main result is presented and proved in Section 4. Section 5 deals with applications of our results to measure dynamical systems. A systematic way to generate such systems from subshifts is discussed in Section 6. The final Section 7 is then devoted to a discussion of singular continuous spectrum for certain models.

1. Measure perturbations and the oracle theorem. We follow [3] and say that a signed Radon measure μ on \mathbb{R} is *translation bounded* if there is an open set U with $\sup_{t \in \mathbb{R}} |\mu|(U+t) < \infty$; we write \mathcal{M} for the space of all signed Radon measures and \mathcal{M}^∞ for the subspace of all translation bounded measures. Here $|\mu|$ denotes the variation of μ . It is easy to see that

$$\mathcal{M}^\infty = \bigcup_{C \geq 0} \mathcal{V}^C,$$

where the latter is defined as in [31], i.e.,

$$\mathcal{V}^C = \{\mu \in \mathcal{M} \mid |\mu|(J) \leq C \max\{|J|, 1\} \text{ for all compact intervals } J \subset \mathbb{R}\}.$$

As explained in [31], there is a metric d on \mathcal{V}^C that will be used later on. The shift is denoted differently here: we let $T_x \mu(A) := \mu(A-x)$ so that $\text{supp } T_x \mu = \text{supp } \mu + x$; in particular, $T_x = S_{-x}$ for the shift S_x used in [31]. In the latter article, the author relies upon [1], as far as the definition of $-\Delta + \mu$ is concerned. Here we sketch a more traditional way via quadratic forms that is pretty easy and direct, since translation bounded measures are form small with respect to $-\Delta$ in one dimension. Here, of course, the form h_0 associated to $-\Delta$ is defined on $W^{1,2}(\mathbb{R})$ by

$$h_0(u) := \int_{\mathbb{R}} u'(t) \bar{u}'(t) dt.$$

Lemma 1.1. *Let $\mu \in \mathcal{M}^\infty$. Then μ is form small with respect to $-\Delta$. In particular,*

$$D(h_0 + \mu) := W^{1,2}(\mathbb{R}), (h_0 + \mu)[u, v] := \int_{\mathbb{R}} u'(t)\bar{v}'(t)dt + \int_{\mathbb{R}} u(t)\bar{v}(t)d\mu(t)$$

defines a closed semibounded form. The associated selfadjoint operator is denoted by $-\Delta + \mu$.

Let us *sketch the proof*:

We use the Sobolev embedding theorem that implies that every $u \in W^{1,2}(\mathbb{R})$ has a continuous representative and use this representative in

$$\mu[u, v] := \int_{\mathbb{R}} u(t)\bar{v}(t)d\mu(t).$$

That the so defined form is form small follows from the Sobolev embedding theorem as well. See [18], Lemma 8 for an appropriate version that gives, for any $\gamma > 0$ a C_γ such that

$$\begin{aligned} |\mu[u, u]| &= \int_{\mathbb{R}} |u(t)|^2 d\mu(t) \\ &= \sum_{k \in \mathbb{Z}} \int_{(k, k+1]} |u(t)|^2 d\mu(t) \\ &\leq \sum_{k \in \mathbb{Z}} \|u|_{[k, k+1]}\|_\infty^2 \\ &\leq \sum_{k \in \mathbb{Z}} (\gamma \|u'|_{[k, k+1]}\|_2^2 + C_\gamma \|u|_{[k, k+1]}\|_2^2) \\ &= \gamma h_0[u, u] + C_\gamma \|u\|_2^2. \end{aligned}$$

That form small perturbations of closed forms yield closed forms is a fundamental truth, see the KLMN-Theorem in [29] or [11], Part I, Paragraph 5. This finishes the sketch of the proof.

Remark 1. (1) The form method can be used in any dimension; see [36, 37] for details and references.

(2) In [1] an alternative direct definition of Schrödinger operators with measure perturbations is given. It has the advantage that the negative part of the measure need not be relatively bounded with respect to the Laplacian, so that operators can be treated that are unbounded from below.

(3) That these two different methods lead to the same operators is not explicitly stated in [1]. It follows, however, from the reasoning in the proof of Theorem 6.1 in [31] (and is proved in [32]).

We now state Remlings Oracle Theorem from [31] for the readers convenience:

Theorem 1.2 (Oracle Theorem). *Let $A \subset \mathbb{R}$ be a Borel set of positive measure, and let $\varepsilon > 0$, $a, b \in \mathbb{R}$, $a < b$, $C > 0$. Then there exist $L > 0$ and a continuous function (the oracle)*

$$\Delta : \mathcal{V}_{(-L, 0)}^C \rightarrow \mathcal{V}_{(a, b)}^C$$

so that the following holds. If $\mu \in \mathcal{V}^C$ and the half line operator $H = -\Delta + \mu \upharpoonright_{[0, \infty)}$ satisfies $\Sigma_{ac}(H) \supset A$, then there exists an $x_0 > 0$ so that for all $x \geq x_0$, we have

$$d(\Delta(1_{(-L, 0)} S_x \mu), 1_{(a, b)} S_x \mu) < \varepsilon.$$

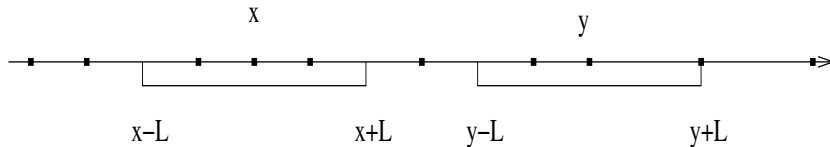


FIGURE 1. A Delone set of finite local complexity

Here, $\Sigma_{ac}(H)$ is the set of energies, for which the density of the absolutely continuous spectral measure is positive.

2. Notions of finite local complexity for sets. We describe the positions of ions in terms of a discrete set $D \subset \mathbb{R}^d$. In case of a perfect crystal, D would be *periodic*, i.e., invariant under translations by elements of a lattice. Here lattice means a discrete subgroup that spans \mathbb{R}^d as an \mathbb{R} -vector space. Recall that a *Delone set* $D \subset \mathbb{R}^d$ is characterized by two properties that it shares with periodic sets: It is *uniformly discrete* and *relatively dense* i.e. satisfies

- there is $r > 0$ such that $U_r(x) \cap U_r(y) = \emptyset$ for $x, y \in D, x \neq y$,
- and there exists $R > 0$ such that $\bigcup_{x \in D} B_R(x) = \mathbb{R}^d$.

Here, the open and closed balls of radius $s > 0$ around $x \in \mathbb{R}^d$ are denoted by $U_s(x)$ and $B_s(x)$ respectively. Delone sets may be used to describe the positions of ions in solids that are not as ordered as crystals. The following property clearly holds true for periodic sets:

A discrete set $A \subset \mathbb{R}^d$ is said to be of *finite local complexity*, provided the following holds: for any $L \geq 0$

$$\{B_L(x) \cap (A - x) \mid x \in A\}$$

is a finite set of subsets of \mathbb{R}^d . We are mainly concerned with subsets of the real line in which case the property means that in any interval $[x - L, x + L]$, $x \in A$ only finitely many different patterns occur. We do not go into formalizing the concept of pattern here, see [24] for a thorough discussion and Figure 1 for an illustration.

Since the geometry of \mathbb{R} is particularly simple, we can easily verify the following observation. (For investigation of higher dimensional analogues to these facts, we refer to [19].)

Proposition 1. (a) *A closed set $D \subset \mathbb{R}$ is a Delone set if it can be written as $D = \{x_n \mid n \in \mathbb{Z}\}$ where $x_{n+1} - x_n \in [2r, \frac{R}{2}]$; it is then of finite local complexity if, additionally, $x_{n+1} - x_n$ can only take finitely many values.*
 (b) *A set of finite local complexity is uniformly discrete but not necessarily relatively dense.*

Given this observation, it is different whether one bases a concept of finite local complexity on “what can be seen around a given point” or on “what comes next”. This is further illustrated by the following examples.

Examples.

- (1) Any periodic set is a Delone set of finite local complexity.
- (2) Any subset of a set of finite local complexity has finite local complexity, e.g., any subset of \mathbb{Z} is of finite local complexity.

It will be convenient to deal with colored Delone sets as well. Let A be a finite set. A colored Delone set with colors from A is a pair (D, a) consisting of a Delone set D and a map

$$a : D \longrightarrow A.$$

The restriction of a colored Delone set $D' = (D, a)$ to a subset B is defined to be the pair

$$D' \cap B := (D \cap B, a : D \cap B \longrightarrow A).$$

A colored Delone set is said to have finite local complexity if the underlying Delone set has finite local complexity. As there are only finitely many elements in A this can easily be seen to be equivalent to finiteness of the set of restriction

$$\{(D' - x) \cap B_R : x \in D\}$$

for each $R > 0$.

3. Notions of finite local complexity for measures. The kind of potentials we are interested in are of the form

$$V_D(t) = \sum_{x \in D} V(t - x),$$

where $V \in L^1$ is a given single site potential of compact support and D is uniformly discrete. A typical piece of such a potential is depicted in Figure 2. More generally, one may be interested in potentials of the form

$$V_{(D,a)} := \sum_{x \in D} V_{a(x)}(t - x)$$

where (D, a) is a Delone colored by the finite A and $V_a, a \in A$, are L^1 potentials with compact support. From the analytical point of view, it does not make much difference to include finite signed measures of compact support as single site potentials;

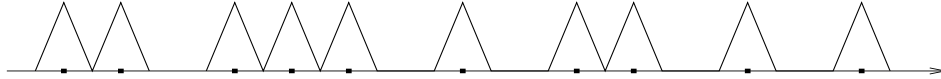


FIGURE 2. A typical potential V_D

If we do not want to use the construction of V_D , how can we formalize the finite local complexity it inherits from D ? The problem we encounter here is the fact that although from a certain perspective, V_D is made of a finite number of parts, we have to pick the right parts. If we slice the real line into intervals in different ways, we end up with an infinite supply of different forms. For point sets as above, this difficulty does not occur, as the point set itself gives a natural grid. Note that point sets can naturally be interpreted as measures: any uniformly discrete D induces the translation bounded measure

$$\delta_D := \sum_{x \in D} \delta_x.$$

Then V_D defined above is just given as $\delta_D * V$, where $*$ denotes the usual convolution. Similarly, we may define a convolution between Delone sets colored by A and tuples of functions $V_a, a \in A$, by

$$(D, a) * V := V_{(D,a)}.$$

We now present a number of different concepts one might come up with. An ultimate criterion of its usefulness for our present purpose lies in the applicability to spectral theory we have in mind, especially whether it can be used together with the oracle theorem, mentioned above. In our context, it will turn out that the simple finite decomposition property s.f.d.p. given in Definition 3.3 gives what we need for our purposes.

We start with some general notation used in our dealing with measures and their local restrictions: A *piece* is a pair (ν, I) consisting of a closed interval I with length $|I| > 0$ and a measure ν supported on I ; we write ν^I as shorthand notation, or sometimes only ν , which is, of course, an abuse of notation. If μ is a measure on \mathbb{R} we say that ν^I *occurs in μ at x* if $(1_{[x, x+|I|]}\mu, [x, x+|I|])$ is a translate of (ν, I) (in the obvious sense). We call $|I|$ the length of the piece ν^I . A *finite piece* is a piece with finite length. The *concatenation* $\nu^I = \nu_1^{I_1} | \nu_2^{I_2} | \dots$ of a finite or countable family $(\nu_j^{I_j})_{j \in N}$, $N = \{1, 2, \dots, |N|\}$ of finite pieces is easy to visualize and somewhat complicated to write down:

$$I = [\min I_1, \min I_1 + \sum_{j \in N} |I_j|]$$

$$\nu = \nu_1 + \sum_{k \in N, k \geq 2} T_{(\min I_1 + \sum_{l=1}^{k-1} |I_l| - \min I_k)} \nu_k.$$

In this case we also say that ν^I is *decomposed into* $(\nu_j^{I_j})_{j \in N}$. Then, $\nu_j^{I_j}$ is the measure that occurs at x for

$$x \in [\min I_j + \sum_{k < j} |I_k|, \max I_j + \sum_{j < k} |I_k|].$$

See Figure 3 for an illustration. What the formulae express is easy to describe: we

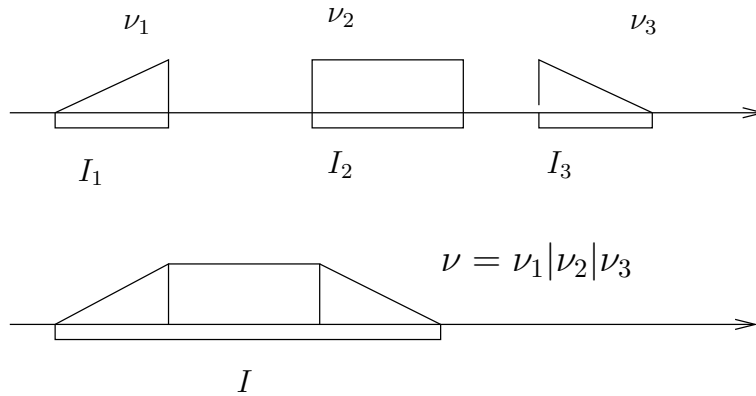


FIGURE 3. The concatenation of pieces.

translate the piece $\nu_2^{I_2}$ in such a way that it starts where $\nu_1^{I_1}$ ends. Then we pick the next piece and put it right to the right of $\nu_2^{I_2}$, and so on.

Now, here comes one version of finite local complexity for measures.

Definition 3.1. A signed measure μ has a *finite number of local pieces* with parameter $\rho > 0$, abbreviated *f.l.p* if for every $x \in \mathbb{R}$ there is $x' \in \mathbb{R}$ with $|x - x'| \leq \rho$ such that, for every $L \geq 2\rho$,

$$\{1_{[-L,L]}(T_{x'}\mu) \mid x \in \mathbb{R}\}$$

is a finite set of measures. We say that μ is *finally f.l.p*, if there is $x_0 \in \mathbb{R}$ such that $1_{[x_0,\infty)}\mu$ has f.l.p.

- Remark 2.** (1) If D is a Delone set of finite local complexity, then δ_D has f.l.p. In fact, we can pick $\rho = R$ and $x' \in D \cap [x - \rho, x + \rho]$, the latter set being nonempty due to the relative denseness condition.
- (2) If D is any subset of a periodic set, then δ_D has f.l.p. In particular, there are f.l.p point measures such that the underlying sets are no Delone sets.

Another version of finite local complexity is given as follows.

Definition 3.2. A measure μ has the *finite decomposition property*, abbreviated f.d.p, if there exists a finite set \mathcal{P} of finite pieces, called local pieces, and $x_0 \in \mathbb{R}$ such that $1_{[x_0,\infty)}\mu^{[x_0,\infty)}$ is the translate of a concatenation $\nu_1^{I_1}|\nu_2^{I_2}|\dots$ with $\nu_j^{I_j} \in \mathcal{P}$, for all $j \in \mathbb{N}$. We denote $x_1 := x_0 + |I_1|, x_2 := x_1 + |I_1|, \dots$ and call the set $\{x_0, x_1, \dots\}$ the *grid* of the decomposition. Without restriction we can and will assume that $\min I = 0$ for every $\nu^I \in \mathcal{P}$. We denote by $l_{\mathcal{P}}, L_{\mathcal{P}}$ the minimal and maximal length of pieces in \mathcal{P} , respectively.

For our main result we will need a stronger version of f.d.p.

Definition 3.3. A μ has the *simple finite decomposition property*, abbreviated s.f.d.p, if it has the f.d.p with a decomposition such that there is $\ell > 0$ with the following property: Assume that the two pieces

$$\nu_{-m}^{I_{-m}}|\dots|\nu_0^{I_0}|\nu_1^{I_1}|\nu_2^{I_2}|\dots|\nu_{m_1}^{I_{m_1}} \quad \text{and} \quad \nu_{-m}^{I_{-m}}|\dots|\nu_0^{I_0}|\mu_1^{J_1}|\mu_2^{J_2}|\dots|\mu_{m_2}^{J_{m_2}}$$

occur in the decomposition of μ with a common first part $\nu_{-m}^{I_{-m}}|\dots|\nu_0^{I_0}$ of length at least ℓ and such that

$$1_{[0,\ell)}(\nu_1^{I_1}|\nu_2^{I_2}|\dots|\nu_{m_1}^{I_{m_1}}) = 1_{[0,\ell)}(\mu_1^{J_1}|\mu_2^{J_2}|\dots|\mu_{m_2}^{J_{m_2}}),$$

where the $\nu_j^{I_j}, \mu_k^{J_k}$ are from the decomposition (in particular, all the pieces belong to \mathcal{P} and so they start at 0) and the latter two concatenations are of lengths at least ℓ . Then

$$\nu_1^{I_1} = \mu_1^{J_1}.$$

Note, that instead of the equality of the pieces we ask for in the last equation of the definition it suffices to ask for the equality of the length of the pieces.

Before we illustrate the preceding definition with examples let us explain the difficulties that might occur and prevent a measure with f.d.p from satisfying s.f.d.p. Essentially there are two problems viz periodicity and the other is local constancy.

Let us discuss periodicity first: Among our local pieces there may be ν and ν' with ν' equal to a concatenation of finitely many copies of ν . Then, a high power of ν can not be decomposed uniquely. This problem may be circumvented by deleting ν' from our set of local pieces.

We now turn to local constancy. You might decompose a fixed multiple of Lebesgue measure - e.g. the measure zero - in pieces of two different lengths, say, in such a way that it does not have the s.f.d.p (whereas f.d.p holds): take one short

piece and then a long one, then two short pieces and then a long one, etc. For a decomposition of that fashion, you can never predict the appearance of a long piece. Note that this situation is not covered by the periodic situation (as the lengths of the pieces may be rationally independent). In some sense these are the only ways how a measure with f.d.p can fail to satisfy s.f.d.p:

Proposition 2. *Let μ be a measure with f.d.p. with respect to the local pieces $\nu_1^{I_1}, \dots, \nu_N^{I_N}$. If μ does not have s.f.d.p. with respect to any set of local pieces, then there must be two local pieces among the $\nu_j^{I_j}$, which are multiples of Lebesgue measure i.e. have the form $c \mathbb{1}_I dx$ for suitable constants c and intervals I .*

Proof. Denote by D_j the set of occurrences of $\nu_j^{I_j}$ in μ . (Here, we count all occurrences irrespective whether these occurrences match with the given decomposition.) Then, $D := \cup D_j$ is relatively dense, as there exists a decomposition of μ . We consider two cases.

D is uniformly discrete. In this case, we can just recode μ in the following way: Let

$$\dots x_{-1} < x_0 < x_1 < \dots$$

be the points of D . Then, we can decompose μ according to

$$\dots \mu_{[x_{-1}, x_0)} | \mu_{[x_0, x_1)} | \mu_{[x_1, x_2)} | \dots$$

As μ has f.d.p with respect to the $\nu_j^{I_j}$, this is indeed a decomposition into finitely many different pieces. By construction the occurrence of each piece is determined locally i.e. only depends on the restriction of μ to a certain neighborhood of the piece. This gives the desired simplicity.

D is not uniformly discrete. Then, there must exist an accumulation point among the occurrences of the local pieces. If the occurrences of a fixed local piece accumulate, then this piece must be locally translation invariant and hence a multiple of the Lebesgue measure. If there are two local pieces with accumulation points we therefore obtain two pieces which are multiples of Lebesgue measure.

Thus, it remains to show that if there is only one local piece with accumulation points, then the system still has s.f.d.p.: In this case, there is only one exceptional piece which is a multiple of Lebesgue measure. If this is the only piece in the decomposition, s.f.d.p is obviously valid. Otherwise, we can use a coding by the occurrences of the other local pieces and break gaps resulting from multiple occurrences of the exceptional piece by starting at the leftmost occurrence of the piece in the gap and then proceeding by adding this piece step by step as long as possible. By f.d.p. this will result in only finitely many pieces. \square

Remark 3. Let us note that even in non-periodic situations it can be hard to decompose into a fixed set of pieces. This is well known even for Delone sets. For example, for certain Delone sets coming from so called primitive substitutions uniqueness of some natural decompositions has been investigated (and proven) in [35] (see [28] as well for the corresponding result in symbolic dynamics).

Examples.

- (1) If μ is finally periodic, i.e., there exist $x_0 \in \mathbb{R}$ and $p > 0$ such that $\mu(A + p) = \mu(A)$ for all measurable $A \subset [x_0, \infty)$, then μ has the s.f.d.p. In fact, $\mathcal{P} := \{(1_{[x_0, x_0+p)}\mu, [x_0, x_0 + p])\}$ does the job, as one easily verifies.

- (2) If D is a Delone set of finite local complexity, then δ_D has the s.f.d.p. This follows from the next lemma below plus the fact that such measures have the f.e.p (defined in the next definition). We can also define a suitable set \mathcal{P} directly: as noted in Proposition 1, there are only finitely many different distances between neighboring points of D ; denote them by $0 < l_1 < \dots < l_n$. We put $I_j = [0, l_j]$, $\nu_j = \delta_o$ and $\mathcal{P} = \{\nu_j^{I_j} \mid j = 1, \dots, N\}$. Now, for every $x_0 \in D$ we find a decomposition; due to the construction of \mathcal{P} , every decomposition is simple.
- (3) The measure μ can be shown to have f.d.p. if and only if there exists a colored Delone D' set and an A tuple (ν_a) , $a \in A$, of measures with compact support such that $\mu = D' * \nu$.

While the preceding notions were based on “what can be seen around a given point”, the next notion formalizes the “what comes next” point of view. This notion is stronger than the previous ones:

Definition 3.4. A signed measure μ is said to have the *finite extension property*, short *f.e.p* if there exist $\rho > 0$, L_0 and, for any $L \geq L_0$ a finite set \mathcal{E}^L of pieces with interval $[0, L]$ with the following property: for $r \geq \rho$ and any piece $\nu^{[0, r]}$ there exists $r' \in [r - \rho, r]$ such that whenever $\nu^{[0, r]}$ occurs in μ at x , there is a $\hat{\nu}^{[0, L]} \in \mathcal{E}^L$ such that $(1_{[0, r']}\nu + T_{r'}\hat{\nu}, [0, r' + L])$ occurs in μ at x . We say that μ is *finally f.e.p* if there is an x_0 such that the above property is valid for $x \geq x_0$.

An illustration concerning the previous definition is given in Figure 4.

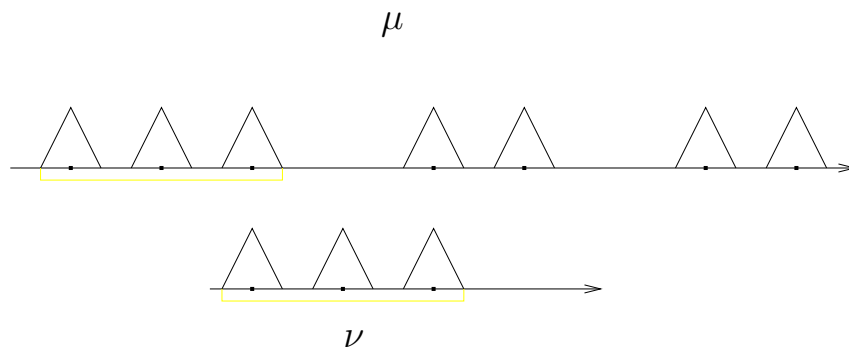


FIGURE 4. A piece ν^I and its occurrence in μ .

- Remark 4.** (1) There is a difference in how we can define finally f.l.p and finally f.e.p: the measure $1_{[x_0, \infty)}\mu$ is almost never f.e.p. In fact, only if it is zero. Otherwise, there are arbitrarily long intervals where μ is zero, followed by something nontrivial. On the other hand, if μ is f.l.p, then $1_{[x_0, \infty)}\mu$ has f.l.p as well
- (2) If D is a Delone set of finite local complexity, then δ_D has f.e.p. In fact, we can pick $\rho = R$ and $x' \in D \cap [x - 2\rho, x]$, the latter set being nonempty due to the relative denseness condition.
- (3) If D is any subset of a periodic set, then δ_D has f.l.p but not necessarily f.e.p.

Finally, we define a strong form of decomposition property.

Definition 3.5. A measure μ has the *unique decomposition property*, abbreviated u.d.p, if it has the finite decomposition property and there exists an $R > 0$ such that the local piece occurring at an x in the decomposition is uniquely determined by the restriction of μ to $[x - R, x + R]$.

The following lemma gathers simple relations between the various notions of local finiteness introduced so far. It is easy to prove.

Lemma 3.6. $u.d.p \implies f.e.p. \implies s.f.d.p \implies f.d.p \implies f.l.p.$

The converse implications do not hold in general: For example Lebesgue measure may easily be decomposed periodically yielding s.f.d.p, still it does not satisfy u.d.p.. Also, we have already discussed how measures with f.d.p may not have s.f.d.p. Finally, considering measures coming from point sets with increasing gaps with incommensurate lengths, one can construct examples of measures with f.l.p. but without f.d.p. However, assuming relative denseness of local pieces we can use 'Deloneification' to show that f.l.p implies u.d.p.. This is the content of the next theorem. The proof is similar to the proof of Proposition 2 above.

Theorem 3.7. *Let μ be a measure on \mathbb{R} with f.l.p. with parameter $\rho > 0$. If for some $L \geq 2\rho$ the set of occurrences of any local piece of length L is relatively dense, then μ is either a multiple of Lebesgue measure or has the unique decomposition property.*

Proof. We consider two cases.

Case 1: There is a local measure ν whose occurrences have a minimal distance $r > 0$. In this case the occurrences of ν in μ form a Delone set. We can use this Delone set to do a recoding of μ by occurrences of ν . More precisely, label the occurrences of ν in μ by

$$\dots, x_{-1} < x_0 < x_1 < x_2 < x_3 \dots$$

Then, we can decompose μ according to

$$\dots \mu|_{[x_{-1}, x_0)} \mu|_{[x_0, x_1)} \mu|_{[x_1, x_2)} \dots$$

As μ has f.l.p and there is a minimal distance between occurrences of ν the Delone set has finite local complexity and this is indeed a decomposition into finitely many different pieces. By construction the occurrence of each piece is determined locally i.e. the unique decomposition property holds.

Case 2: The set of occurrences of each local measure ν has arbitrary small gaps. The assumption means that for each local measure ν

$$\nu = T_x \nu|_{\text{supp } \nu}$$

for a set of nonzero x with $x \rightarrow 0$. Then, ν is (locally) translation invariant and hence is a multiple of Lebesgue measure. Different local measures must be the same multiple as $L \geq 2\rho$. \square

Corollary 1. *Let $S > 0$ be arbitrary. Let μ be a measure on \mathbb{R} with f.l.p. with parameter ρ such that for some $L \geq 2\rho + S$ the set of occurrences of any local piece of length L is relatively dense. Let ν be a measure of compact support contained in $[0, S]$. Then, $\mu * \nu$ is either Lebesgue measure or has unique decomposition property.*

Proof. As μ has f.l.p and ν has compact support, the measure $\mu * \nu$ has f.l.p. as well. By assumption the measure $\mu * \nu$ has relative dense set of occurrences for local measures of length $L - S$. Now, the statement follows from the previous theorem. \square

If the underlying measure μ comes from a Delone set of finite local complexity a somewhat stronger statement holds.

Proposition 3. *Assume that D is a Delone set of finite local complexity and ν is a finite signed measure of compact support. Then*

$$\mu = \sum_{x \in D} T_x \nu = \nu * \delta_D$$

has the s.f.d.p.

Unfortunately, the proof is not as evident as the statement may seem. Here is the problem: You cannot recover D from $\nu * \delta_D$ in general. Easy periodic examples will convince you of this unpleasant fact. This lead to additional restrictions in [15] concerning the single site bumps that were allowed. Here, we can circumvent these additional assumptions, because we can still choose the particular decomposition according to our purposes.

Proof. By shifting D and ν we may suppose that $\min(\text{supp } \nu) = 0$; moreover, $\max(\text{supp } \nu) =: s > 0$ without restriction. Let $D = \{x_k \mid k \in \mathbb{Z}\}$ with the x_k strictly increasing. We decompose μ starting from x_0 and take $\{x_k \mid k \in \mathbb{Z}, k \geq 0\}$ as a grid. Note that

$$1_{[x_k, x_{k+1}]} \mu = \sum_{x \in D \cap [x_k - s, x_k]} T_x \nu = T_{x_k} \sum_{x \in D \cap [x_k - s, x_k]} T_{x - x_k} \nu,$$

with a finite number of different possibilities for $x - x_k$ in $[-s, 0]$, independently of x_k . Therefore,

$$\mathcal{P} := \left\{ \left(\sum_{x \in D \cap [x_k - s, x_k]} T_{x - x_k} \nu, [0, x_{k+1} - x_k] \right) \mid k \in \mathbb{Z}, k \geq 0 \right\}$$

is a finite set of finite pieces that decomposes $1_{[x_0, \infty)} \mu$. Now we want to show that the decomposition we chose has the s.f.d.p. Choose $l > \max\{s, L_{\mathcal{P}}\}$ and assume that two pieces $\nu_{-m}^{I_{-m}} | \dots | \nu_0^{I_0} | \nu_1^{I_1} | \nu_2^{I_2} | \dots | \nu_{m_1}^{I_{m_1}}$ and $\nu_{-m}^{J_{-m}} | \dots | \nu_0^{J_0} | \mu_1^{J_1} | \mu_2^{J_2} | \dots | \mu_{m_2}^{J_{m_2}}$ occur in the decomposition of μ .

Denote the point in the grid where $\nu_1^{I_1}$ starts by x and the point where $\mu_1^{J_1}$ starts by x' . Since μ is given as a sum of translates depending on what is happening to the left, we get that

$$1_{[x, x + \ell_{\mathcal{P}}]} \mu = T_{x' - x} 1_{[x', x' + \ell_{\mathcal{P}}]} \mu.$$

Subtracting the common contribution from $D \cap [x - s, x]$ and $D \cap [x' - s, x']$ we get

$$T_{|I_1|} \nu = T_{|J_1|} \nu,$$

and the asserted equality of the length $|I_1| = |J_1|$ follows. \square

The considerations of this section suggest the following as definition for a Delone measure of finite local complexity.

Definition 3.8. A measure on \mathbb{R} is said to be a Delone measure of finite local complexity if there exist finitely many measures ν_1, \dots, ν_N with compact support such that with the sets D_j of occurrences of ν_j in μ , $j = 1, \dots, N$, the following holds:

- The set $\cup_j D_j$ is a Delone set of finite local complexity.
- For any $x \in \text{supp } \mu$, there exists a $j \in \{1, \dots, N\}$ and an $p \in D_j$ such that x belongs to $p + \text{supp } \nu_j$ and μ agrees with $T_p \nu_j$ on $p + \text{supp } \nu_j$.

Let us note that this definition can be used verbatim to deal with measures in arbitrary dimensions.

4. The main result. In this section, we present the main result on absence of absolutely continuous spectrum for aperiodic measure of 'finite local complexity'.

Theorem 4.1. *Let $\mu \in \mathcal{M}^\infty$ be a measure that has the s.f.d.p and assume that the absolutely continuous spectrum of the half-line operator $-\Delta + \mu \upharpoonright_{[0, \infty)}$ is not empty. Then, μ is eventually periodic.*

Proof. The idea goes as follows: due to the s.f.d.p there are pieces of arbitrary length that appear infinitely often. If we pick such a piece of sufficient length, the oracle theorem allows us to predict the next piece up to a small error. Since the next piece itself comes from a finite set, it is uniquely determined, provided the error is small enough. That means that after the two occurrences of the given long piece, the same shorter piece follows. Iterating the procedure we get longer and longer translates of the measure that agree, which means that the measure is eventually periodic.

Here are the details: Assume that $\sigma_{ac}(-\Delta + \mu) \neq \emptyset$ is nonempty. Then, $A = \Sigma_{ac}(-\Delta + \mu)$ has positive measure. Let $G := \{x_k \mid k \in \mathbb{N}_0\}$ be the grid and \mathcal{P} the set of pieces of a decomposition with the s.f.d.p and ℓ the according length.

Denote

$$\mathcal{D}_\ell := \{1_{(-1, \ell)}[T_{-x}\mu] \mid x \in G\}$$

and note that this set is finite, due to the s.f.d.p. Consequently, there is $\varepsilon > 0$ with $d(\nu_1, \nu_2) > 2\varepsilon$ for different $\nu_1, \nu_2 \in \mathcal{D}_\ell$.

This fixes ε for our application of the oracle theorem; moreover, we choose $a = -1, b = \ell$. Pick L according to the oracle theorem, without restriction $L > \ell$.

We construct a coarser grid $G_L \subset G$ by $y_0 := x_0$ and $y_{k+1} - y_k \in [L, L + L\mathcal{P}]$. Since $G \cap [y_k, y_{k+1}]$ is finite, it is clear that

$$\mathcal{P}_L := \{(T_{-y_k}[1_{[y_k, y_{k+1}}]\mu], T_{-y_k}([y_k, y_{k+1}] \cap G)) \mid k \in \mathbb{N}\}$$

is finite. Consequently, there is at least one k , for which infinitely many translates of $1_{[y_k, y_{k+1}}]\mu$ appear in μ so that the corresponding parts of G are translates of $G \cap [y_k, y_{k+1}]$. Call $z^1 := y_{m+1}$, where y_m is one of the corresponding points in G_L , $m > k$ and denote $y^1 := y_{k+1}$; see the following Figure 5. We apply the oracle

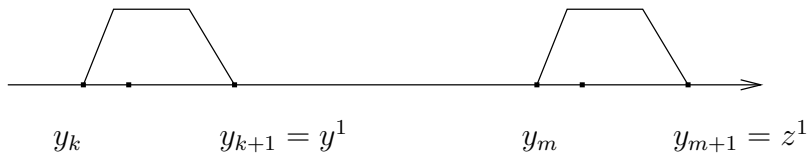


FIGURE 5. Illustration of the proof.

theorem with $x = y^1$ and $x = z^1$ and get that

$$d(1_{(-1,\ell)}T_{y^1}\mu, 1_{(-1,\ell)}T_{z^1}\mu) \leq 2\varepsilon.$$

By our choice of ε this implies

$$1_{(-1,\ell)}T_{y^1}\mu = 1_{(-1,\ell)}T_{z^1}\mu.$$

By definition of \mathcal{P}_L we moreover know that the pieces of μ starting at y_k and y_m , respectively, are decomposed in the same way. The s.f.d.p (note that $L > \ell$) guarantees that the pieces start at y^1 and z^1 are translates of each other, so that the next point $y^2 := \min(G \cap (y^1, \infty))$ in G following y^1 and $z^2 := \min(G \cap (z^1, \infty))$ satisfy

$$z^2 - z^1 = y^2 - y^1 \text{ and } T_{-y^1}[1_{[y^1, y^2]}\mu] = T_{-z^1}[1_{[z^1, z^2]}\mu].$$

This implies that $1_{[y_k, y^2]}\mu$ is a translate of $1_{[y_m, z^2]}\mu$. We can now iterate the procedure to find $y^3, z^3 \in G$ such that $1_{[y_k, y^3]}\mu$ is a translate of $1_{[y_m, z^3]}\mu$. Since, moreover, $y^{j+1} - y^j, z^{j+1} - z^j \geq \ell_{\mathcal{P}}$, we get that $1_{[y_k, \infty)}\mu$ is a translate of $1_{[y_m, \infty)}\mu$. In other words, μ is finally periodic. \square

5. Application to measure dynamical systems. In this section we consider applications of the main result to dynamical systems of measures.

A dynamical system (Ω, S) consists of a compact set Ω and a continuous action S of \mathbb{R} on Ω . A dynamical system is called minimal if

$$\{S_x\mu : x \in \mathbb{R}\}$$

is dense in Ω for every $\mu \in \Omega$. An S invariant probability measure on a dynamical system (Ω, S) is called ergodic if any invariant measurable subset of Ω has measure zero or measure one. If there exists a unique invariant probability measure, then (Ω, S) is called uniquely ergodic. An element $\omega \in \Omega$ is called periodic with period $x \neq 0$ if $S_x\omega = \omega$. The dynamical system is called periodic if all of its elements are periodic with the same period.

We will be concerned with measure dynamical systems. In this case Ω is a closed (and hence compact) subset of \mathcal{V}^C for some $C > 0$, which is invariant under the shifts $S = T$. A measure dynamical system is said to have one of properties defined in the previous section if each element of Ω has this property. Each measure dynamical system (Ω, T) gives rise to the family $H_\omega, \omega \in \Omega$, defined by

$$H_\omega = -\Delta + \omega.$$

The first main result in this section concerns minimal measure dynamical systems with a finite local complexity property.

Theorem 5.1. *Let (Ω, T) be a minimal measure dynamical system with s.f.d.p. Then, (Ω, T) is periodic if one H_ω possesses some absolutely continuous spectrum.*

Proof. Let $\omega \in \Omega$ be chosen such that H_ω possesses some absolutely continuous spectrum. Then, the restriction of H_ω to $[0, \infty)$ or the restriction of H_ω to $(-\infty, 0]$ must possess some absolutely continuous spectrum by general principles. Assume without loss of generality that the restriction of H_ω to $[0, \infty)$ possesses some absolutely continuous spectrum. Then, ω is eventually periodic by our main result, Theorem 4.1. By minimality, then each element of Ω must be periodic (with the same period). \square

Remark 5. If the system is minimal then the absolutely continuous spectrum of the H_ω does not depend on ω by results of Last /Simon [20] (if the measures possess a density with respect to Lebesgue measure).

Our second main result in this section deals with ergodic measure dynamical systems. We need a little preparation.

Lemma 5.2. *Let (Ω, S) be a dynamical system with ergodic measure m . If the set of eventually periodic points has positive measure, then any point in the support of m is periodic (with the same period).*

Proof. For each compact interval I let A^I be the set of points of Ω which are eventually periodic with period $p \in I$. Then, each A^I is measurable and invariant under S and hence has measure zero or one by ergodicity. As A^I and A^J are disjoint for $I \cap J = \emptyset$, we infer from the assumption and an intersection procedure that there exists exactly one p such that $A^p := A^{\{p\}}$ has full measure. Then $S_n A^p$ has full measure for each $n \in \mathbb{N}$ and so has

$$A_\infty^p := \bigcap_{n \in \mathbb{N}} S_n A^p.$$

By construction the elements of A_∞^p are periodic with period p . We have thus found a set of full measure all of whose elements are periodic (with the same period). This implies the statement. \square

Whenever (Ω, T) is a measure dynamical system with ergodic measure μ then there exists a closed subset Σ_{ac} of \mathbb{R} such that Σ_{ac} is the absolutely continuous spectrum of H_ω for μ almost every $\omega \in \Omega$ (see e.g. the monograph [6]).

Theorem 5.3. *Let (Ω, T) be measure dynamical system of finite local complexity and m an ergodic measure on Ω . If $\Sigma_{ac} \neq \emptyset$, then ω is periodic for every ω in the support of m .*

Proof. By our main result Theorem 4.1 the measure ω is eventually periodic for almost every $\omega \in \Omega$. This implies the statement by the previous lemma. \square

6. Systems generated by subshifts. In this section, we discuss a systematic way to generate measure dynamical systems with f.l.p by a suspension type construction over a subshift over a finite alphabet. In the last part of the section we then show how this can be applied to deal with Delone sets.

We start by introducing the necessary notation concerning subshifts. Let A , called the alphabet, be a finite set equipped with the discrete topology. A pair (X, τ) is a subshift over A if Ω is a closed subset of $A^{\mathbb{Z}}$, where $A^{\mathbb{Z}}$ is given the product topology and Ω is invariant under the shift operator $\tau : A^{\mathbb{Z}} \rightarrow A^{\mathbb{Z}}$, $\tau a(n) \equiv a(n+1)$. We consider sequences over A as words and use standard concepts from the theory of words, see e.g. [27]. In particular, the number of occurrences of v in w is denoted by $\#_v(w)$ and the length $|w|$ of the word $w = w(1) \dots w(n)$ is given by n . Given a subshift (X, τ) over the finite alphabet A with cardinality N , we can associate to each $\ell = (l_1, \dots, l_N) \in (0, \infty)^N$ a new measure dynamical system (X_ℓ, T) over \mathbb{R} by a suspension type construction. More precisely, assume without loss of generality that $A = \{1, \dots, N\}$ and define for each $x \in X$ the measure ξ_x by

$$\xi_x := \sum_{j \in \mathbb{Z}} x(j) \delta_{\sum_{k=0}^{j-1} l_{x(k)}}.$$

where δ_x denotes the unit point measure at $s \in \mathbb{R}$. Now, set

$$X_\ell := \{T_t \xi_x : x \in X, t \in \mathbb{R}\}.$$

Then, (X_ℓ, T) is a measure dynamical system, which inherits many properties of (X, τ) . For example it is not hard to see that minimality of (X, τ) implies minimality of (X_ℓ, T) . Similarly, (X_ℓ, T) is uniquely ergodic if (X, τ) is uniquely ergodic. More generally, any invariant measure n on (X, τ) induces a unique measure m on (X_ℓ, T) with

$$n(U_w)|I| = m(U_{w,I})$$

for any finite word w and any sufficiently short interval I . Here,

$$U_w := \{x \in X : w = x(0) \dots x(|w| - 1)\}$$

and

$$U_{w,I} := \{\omega \in X_\ell : w' \text{ occurs in } \omega \text{ at an } t \in I\},$$

with $w' := \sum_{j=1}^{|w|} w(j) \delta_{\sum_{k=0}^{j-1} l_{x(k)}}$. A discussion of essentially this phenomenon can be found in [21, 15].

After these preparations we can now discuss how to associate measure dynamical systems to subshifts. Let (X, τ) be a subshift over the finite A with cardinality N . Let $\nu := (\nu_1, \dots, \nu_N)$ be an N -tuple of measures with compact support such that

$$\inf \text{supp } \nu_j = 0$$

for all $j = 1, \dots, N$. Chose $l_j := \text{supp } \nu_j$ if ν_j is not just a point measure in 0 and $l_j = 1$ otherwise. Now, define for each $x \in X$ the measure ω_x by

$$\omega_x := \sum_{n \in \mathbb{Z}} \delta_{\sum_{k=0}^{n-1} l_{x(k)}} * \nu_{x(n)}.$$

Thus, we replace in ξ_x any point measure with coefficient j with the measure ν_j . Let

$$\Omega_\nu := \{T_t \omega_x : x \in X, t \in \mathbb{R}\}.$$

By construction Ω_ν has *f.d.p* and by the criteria above, in particular, Proposition 2, we have a simple sufficient condition for *s.f.d.p*. Moreover, by construction there is a factor map

$$\pi : (X_\ell, T) \longrightarrow (\Omega_\nu, T)$$

such that $\pi(\xi)$ is the measure arising by replacing any point measure with coefficient j in ξ with the measure ν_j . By standard theory, the factor (Ω_ν, T) inherits properties like minimality and unique ergodicity from (X_ℓ, T) . Moreover, π maps invariant (ergodic) measures on (X_ℓ, T) to invariant (ergodic) measures on (Ω_ν, T) (see e.g. [3] for a discussion of these properties).

Putting the above considerations together we have shown the following.

Proposition 4. *Let (X, τ) be a subshift over the finite alphabet A . Let N be the cardinality of A and $\nu := (\nu_1, \dots, \nu_N)$ be a tuple of measures with compact support and $\inf \nu_j = 0$, $j = 1, \dots, N$. Assume that not two of the ν_j are multiple of Lebesgue measure. Then (Ω_ν, T) is a measure dynamical system with *s.f.d.p* and the following holds:*

- (a) *The measure dynamical system (Ω_ν, T) is minimal if (X, τ) is minimal.*
- (b) *Any invariant probability measure n on (X, τ) induces a canonical invariant probability measure m on (Ω_ν, T) . If n is ergodic, so is m .*
- (c) *If (X, τ) is uniquely ergodic, so is (Ω_ν, T) .*

Combined with our main result, Theorem 4.1, the previous proposition gives the following.

Corollary 2. *Let (X, τ) be a subshift over a finite alphabet with ergodic measure n . Let (Ω_ν, T) with ergodic measure m be associated to (X, τ) as in the previous proposition. If the almost sure absolutely continuous spectrum of the family (H_ω) , $\omega \in \Omega_\nu$, is not empty, then any ω from the support of m is periodic.*

Remark 6. The result does not state periodicity of (X, τ) . In fact, such a periodicity can not be concluded. Take for example all measures ν_j to be equal δ_0 and all lengths l_j equal to 1. Then, (Ω, T) will be periodic irrespective of (X, τ) .

Remark 7. The discussion in this section was phrased in terms of dynamical systems. However, it is clear that we actually have pointwise constructions viz any sequence over A can be turned into a measure with *s.f.d.p* by the above procedure.

Let us here shortly describe two classes of subshifts to which our corollary can be applied:

Example - Bernoulli shifts. Here, the subshift is given by $X = \{0, 1\}^{\mathbb{Z}}$. The measure n on X is the product measure $\Pi := \prod \mu$, where μ is the measure assigning the value p to $\{0\}$ and the value $1 - p$ to $\{1\}$ for some fixed p with $0 < p < 1$. Choosing measures ν_1 and ν_2 with compact support, we can then apply the above theory to conclude absence of absolutely continuous spectrum almost surely whenever the resulting dynamical system is not periodic. Of course, for these more is known and in fact dynamical localization was shown in [10] (at least when the measures in question possess an L^1 density).

Example - Circle maps. Let $\alpha \in (0, 1)$ be irrational and $\beta \in (0, 1)$ be arbitrary. Let $\Omega(\alpha, \beta)$ be the smallest subshift over $\{0, 1\}$ containing the function

$$V_{\alpha, \beta} : \mathbb{Z} \longrightarrow \{0, 1\}, \quad V_{\alpha, \beta}(n) := 1_{(1-\beta, 1]}(n\alpha \pmod{1}).$$

Then, (Ω, α) is aperiodic, minimal and uniquely ergodic as α is irrational (see e.g. Appendix of [8] for discussion). Again, we can conclude absence of absolutely continuous spectrum almost surely whenever the resulting dynamical system is not periodic.

So far, the considerations in this section have been concerned with subshifts and not with Delone sets. Let us finish this section, by showing how Delone sets of finite local complexity give rise to sequences over a finite alphabet. Combined with the above discussion this will provide a way to associate measure dynamical systems to Delone dynamical systems. We only sketch the procedure. Let $\Lambda \subset \mathbb{R}$ be a Delone set of finite local complexity. Thus, for each $R > 0$ the set

$$A := \{(x - R, x + R) \cap \Lambda : x \in \Lambda\}$$

is finite. Assume without loss of generality that $0 \in \Lambda$ and order the points of Λ according to

$$\dots x_{-2} < x_{-1} < 0 = x_0 < x_1 < x_2 \dots$$

We can then associate to Λ the sequence

$$x : \mathbb{Z} \longrightarrow A, \quad x_j \mapsto \{(x_j - R, x_j + R) \cap \Lambda.$$

Obviously, Λ can be recovered from x and vice versa (see [23] for a further discussion).

7. Absence of point spectrum and purely singular continuous spectrum.

In this section, we shortly discuss a criterion for absence of point spectrum. This will allow us to provide some examples with purely singular continuous spectrum. In the discrete case the criterion was introduced by Gordon [12]. A version for continuum operators was given by Damanik/Stolz [9] (see [32] for recent results as well). In the main result of this section our measures will be assumed to have a density with respect to Lebesgue measure. The condition of translation boundedness then amounts to a uniform L^1_{loc} property. Accordingly, we will speak about L^1_{loc} dynamical systems.

We start by discussing a condition first used by Kaminaga [14] in his study of absence of eigenvalues for discrete Schrödinger operators associated to circle maps. The condition has then been used in various versions in the study of discrete one-dimensional quasicrystals (see [7] for a survey with further references). We are going to use it in an analogous way.

Definition 7.1. A measure dynamical system (Ω, T) with ergodic measure m is said to satisfy condition (K) if there exist $P_n \geq 0$ with $P_n \rightarrow \infty$ such that

$$G_n := \{\omega \in \Omega : \omega_{[0, P_n)} = (T_{P_n} \omega)_{[0, P_n)} = (T_{-P_n} \omega)_{[0, P_n)}\}$$

satisfies

$$\limsup_{n \rightarrow \infty} m(G_n) > 0.$$

The relevance of this condition in the study of absence of eigenvalues comes from the so-called Gordon lemma. Here comes a special case of the 'Gordon Lemma' from [9].

Lemma 7.2. *Let V be function such that $V dx$ is a translation bounded measure. If there exist P_n -periodic translation bounded measures of the form $V_n dx$ with $P_n \rightarrow \infty$ and*

$$\int_{-P_n}^{2P_n} |V(x) - V_n(x)| dx = 0$$

for all n , then $-\Delta + V$ does not have eigenvalues.

We are now in a position to state an abstract result on absence of eigenvalues.

Lemma 7.3. *If the L^1_{loc} dynamical system (Ω, T) with ergodic measure m satisfies condition (K), then $-\Delta + \omega$ does not have eigenvalues for m almost every $\omega \in \Omega$.*

Proof. By ergodicity, the set Ω_c of $\omega \in \Omega$ with empty point spectrum has either full measure or zero measure. Thus, to prove the statement it suffices to show that it has positive measure. By the previous Gordon lemma, we have that

$$G := \bigcap_n \bigcup_{k=n}^{\infty} G_k$$

belongs to Ω_c . By (K) and a short calculation we obtain

$$m(\Omega_c) \geq m(G) \geq \lim_{n \rightarrow \infty} m\left(\bigcup_{k=n}^{\infty} G_k\right) \geq \limsup m(G_n) > 0.$$

This finishes the proof. \square

We can combine the previous lemma with our main result to obtain the following.

Theorem 7.4. *Let (Ω, T) be an L^1_{loc} dynamical system with ergodic measure m satisfying s.f.d.p and (K). Then, $-\Delta + \omega$ has almost surely purely singular continuous spectrum.*

For the class of systems generated by subshifts, the condition (K) can be derived whenever the subshift satisfies the analog condition. We say that a subshift (X, τ) satisfies (K), if there exist natural numbers $p_n \rightarrow \infty$ such that

$$G_n^X := \{x \in X : x(-p_n) \dots x(-1) = x(0) \dots x(p_n - 1) = x(p_n) \dots x(2p_n - 1)\}$$

satisfy

$$\limsup_{n \rightarrow \infty} (G_n^X) > 0$$

Proposition 5. *Let (X, τ) be a subshift over a finite alphabet with ergodic measure m , which satisfies (K). Let (Ω_ν, T) be an associated measure dynamical system with measure m . Then (Ω_ν, T) satisfies condition (K).*

Proof. This follows easily from the defining properties of m . \square

Example - Circle maps. The previous proposition can be applied to all circle maps with α whose continued fraction expansion has infinitely many coefficients of value at least 4 [14]. The set of such α 's is a set of full Lebesgue measure.

Acknowledgments. P.S. and D.L. take the opportunity to thank David Damanik for an invitation to a highly enjoyable conference. Useful discussions with Christian Remling on this occasion are also gratefully acknowledged. Partial support from the German Science Foundation (DFG) is gratefully acknowledged. Finally, the authors would like to thank the referees for their careful reading and suggestions.

REFERENCES

- [1] A. Ben Amor and C. Remling, *Direct and inverse spectral theory of one-dimensional Schrödinger operators with measures*, Integr. Equ. Oper. Theory, **52** (2005), 395–417.
- [2] J. Bellissard, B. Iochum, E. Scoppola and D. Testard, *Spectral properties of one-dimensional quasi-crystals*, Commun. Math. Phys., **125** (1989), 527–543.
- [3] M. Baake and D. Lenz, *Deformation of Delone dynamical systems and pure point diffraction*, J. Fourier Anal. Appl., **11** (2005), 125–150.
- [4] M. Baake and R. V. Moody (eds.), “Directions in Mathematical Quasicrystals,” Amer. Math. Soc., Providence, RI, 2000.
- [5] J. Breuer and R. Frank, *Singular spectrum for radial trees*, preprint, Rev. Math. Phys., **21** (2009), 929–945.
- [6] R. Carmona and J. Lacroix, “Spectral Theory of Random Schrödinger Operators,” Probability and Its Applications, Birkhuser Boston, Inc., Boston, MA, 1990.
- [7] D. Damanik, *Gordon-type arguments in the spectral theory of one-dimensional quasicrystals*, in “Directions in Mathematical Quasicrystals,” CRM Monogr. Ser., **13**, Amer. Math. Soc., Providence, RI, (2000), 277–305.
- [8] D. Damanik and D. Lenz, *Half-line eigenfunction estimates and purely singular continuous spectrum of zero Lebesgue measure*, Forum Math., **16** (2004), 109–128.
- [9] D. Damanik and G. Stolz, *A generalization of Gordon’s theorem and applications to quasiperiodic Schrödinger operators*, Electron. J. Differential Equations, **2000**, 8 pp. (electronic)
- [10] D. Damanik, R. Sims and G. Stolz, *Localization for one-dimensional, continuum, Bernoulli-Anderson models*, Duke Math. J., **114** (2002), 59–100.
- [11] W. G. Faris, “Self-adjoint Operators,” Lecture Notes in Mathematics, Vol. **433**, Springer-Verlag, Berlin-New York, 1975.
- [12] A. Gordon, *The point spectrum of the one-dimensional Schrödinger operator*, Uspehi Mat. Nauk, **31** (1976), 257–258.
- [13] C. Janot, “Quasicrystals: A Primer,” Oxford University Press, Oxford, 1992.

- [14] M. Kaminaga, *Absence of point spectrum for a class of discrete Schrödinger operators with quasiperiodic potential*, Forum Math., **8** (1996), 63–69.
- [15] S. Klassert, “Spektraltheoretische Untersuchungen von zufälligen Operatoren auf Delone-Mengen,” Dissertation 2007, <http://archiv.tu-chemnitz.de/pub/2007/0068/index.html>.
- [16] S. Kotani, *Lyapunov indices determine absolutely continuous spectra of stationary random one-dimensional Schrödinger operators*, in “Stochastic Analysis” (Katata/Kyoto, 1982), North-Holland Math. Library, **32**, North-Holland, Amsterdam, (1984), 225–247.
- [17] S. Kotani, *Jacobi matrices with random potentials taking finitely many values*, Rev. Math. Phys., **1** (1989), 129–133.
- [18] P. Kuchment, *Quantum graphs. I. Some basic structures*, Special section on quantum graphs, Waves Random Media, **14** (2004), S107–S128.
- [19] J. Lagarias, *Geometric models for quasicrystals I. Delone sets of finite type*, Discrete Comput. Geom., **21** (1999), 161–191.
- [20] Y. Last and B. Simon, *Eigenfunctions, transfer matrices, and absolutely continuous spectrum of one-dimensional Schrödinger operators*, Invent. Math., **135** (1999), 329–367.
- [21] J.-Y. Lee, R. V. Moody and B. Solomyak, *Pure point dynamical and diffraction spectra*, Ann. Henri Poincaré, **3** (2002), 1003–1018.
- [22] D. Lenz, *Ergodic theory and discrete one-dimensional random Schrödinger operators: Uniform existence of the Lyapunov exponent*, Contemporary Mathematics, **327** (2003), 223–238.
- [23] D. Lenz, *Aperiodic order via dynamical systems: Diffraction theory for sets of finite local complexity*, Contemp. Math., **485**, (Idris Assani, ed), (2009), 91–112.
- [24] D. Lenz and P. Stollmann, *Delone dynamical systems and associated random operators*, in “Proc. OAMP, Constanta 2001” (eds. Combes et al), Theta Foundation, 2003.
- [25] D. Lenz and P. Stollmann, *Generic sets in spaces of measures and generic singular continuous spectrum for Delone Hamiltonians*, Duke Math. J., **131** (2006), 203–217.
- [26] D. Lenz and P. Stollmann, *Generic subsets in spaces of measures and singular continuous spectrum*, in “Mathematical Physics of Quantum Mechanics,” Lecture Notes in Phys., **690**, Springer, Berlin, (2006), 333–341
- [27] M. Lothaire, “Combinatorics on Words,” Encyclopedia of Mathematics and Its Applications, **17**, Addison-Wesley, Reading, Massachusetts, 1983.
- [28] B. Mossé, *Puissances de mots et reconnaissabilité des points fixes d’une substitution*, Theoret. Comput. Sci., **99** (1992), 327–334.
- [29] M. Reed and B. Simon, “Methods of Modern Mathematical Physics. II. Fourier Analysis, Self-Adjointness,” Academic Press [Harcourt Brace Jovanovich, Publishers], New York-London, 1975.
- [30] C. Remling, *The absolutely continuous spectrum of Jacobi matrices*, Annals of Math., to appear, [arXiv:0706.1101](https://arxiv.org/abs/0706.1101).
- [31] C. Remling, *The absolutely continuous spectrum of one-dimensional Schrödinger operators*, Math. Phys. Anal. Geom., **10** (2007), 359–373.
- [32] C. Seifert, in preparation
- [33] M. Senechal, “Quasicrystals and Geometry,” Cambridge University Press, Cambridge, 1995.
- [34] D. Shechtman, I. Blech, D. Gratias and J. W. Cahn, *Metallic phase with long-range orientational order and no translational symmetry*, Phys. Rev. Lett., **53** (1984), 1951–1953.
- [35] B. Solomyak, *Nonperiodicity implies unique composition for self-similar translationally finite tilings*, Discrete Comput. Geom., **20** (1998), 265–279.
- [36] P. Stollmann, *Smooth perturbations of regular Dirichlet forms*, Proc. Amer. Math. Soc., **116** (1992), 747–752.
- [37] P. Stollmann and J. Voigt, *Perturbation of Dirichlet forms by measures*, Potential Anal., **5** (1996), 109–138.
- [38] A. Sütő, *The spectrum of a quasiperiodic Schrödinger operator*, Commun. Math. Phys., **111** (1987), 409–415.
- [39] A. Sütő, *Singular continuous spectrum on a Cantor set of zero Lebesgue measure for the Fibonacci Hamiltonian*, J. Stat. Phys., **56** (1989), 525–531.

Received January 2010; revised September 2010.

E-mail address: Steffen.Klassert@mathematik.tu-chemnitz.de

E-mail address: daniel.lenz@uni-jena.de

E-mail address: P.Stollmann@mathematik.tu-chemnitz.de